

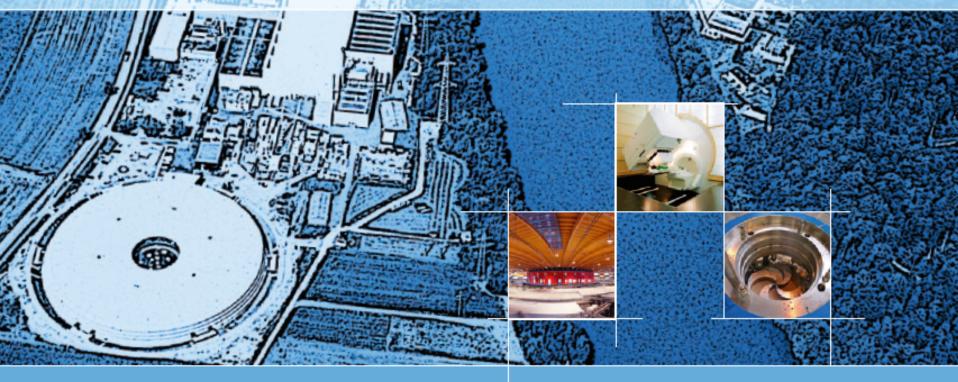
Presented 12/06/2009

Christos Zamantzas



DIPAC 2009

9th European Workshop on Beam Diagnostics and Instrumentation for Particle Accelerators



May 25 – 27, 2009

Basel, Switzerland
Hotel Mercure



Numbers & Facts (unofficial)

2.5 days workshop with

- 10 invited talks [Performance of and Experience with the LHC Beam Diagnostics, R.Jones]
- 14 contributed talks [Digital BPM System for Hadron Accelerators, J. Belleman]
- 3 poster sessions with 115 posters

[14 from BE/BI]

- ~ 200 attendees
- ~ 95 Institutes/Universities
- ~ 7 exhibiting vendors

http://dipac09.web.psi.ch/

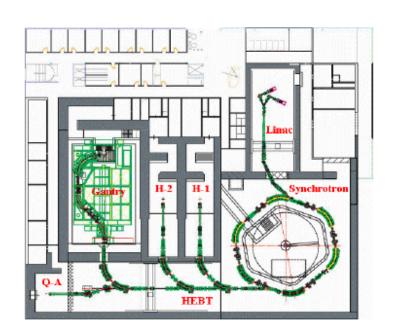


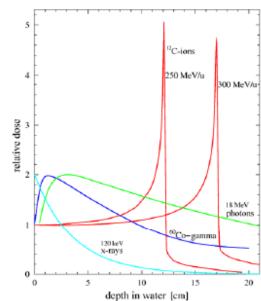
Instrumentation Req. for Dif. Accelerator Types

II. Hadron Therapy Machines

the standards: beam energy, intensity, tunes, orbits not really a problem

Example: parameters of the HIT project in Heidelberg:





control of penetration depth / Bragg peak via beam energy

particles: p, C, He, O

beam energy 50 - 430 MeV/u

beam size 4-10 mm extraction time 1-10 s

extraction intensity 106-4*1010 ions/spill

beam power 360 W dc power



Highlights from the 13th BIW



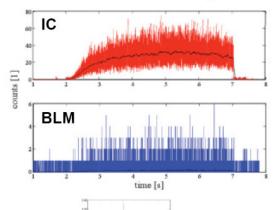
Hadron Therapy Diagnostics

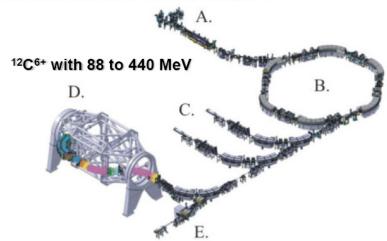


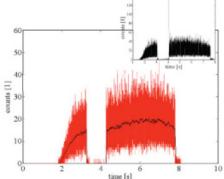
TUPTPF044 — Beam Quality Measurements of the Synchrotron and HEBT of the Heidelberg Ion Therapy Center, T. Hoffmann, D. Ondreka, A. Peters, A. Reiter, M. Schwickert.

Beam diagnostics in hadron therapy facilities plays a very peculiar role.

Reliability, calibration and precision are a "must"!







The "spill structure" is redundantly measured by using ionization chambers, scintillators and beam loss monitors.

The "spill pause" is important to minimize undesired healthy tissue exposure to the beam.

If the spill pause does not work properly a fast "RF knockout" kills the beam within a spill.



Highlights from the 13th BIW



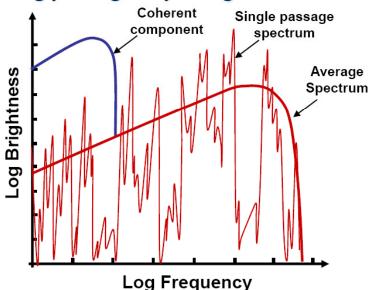
Beam Absolute Bunch Length by Radiation Fluctuation Analysis



THVTIO01 — Recent Beam Measurements and New Instrumentation at the ALS, F. Sannibale

In real beams, due to presence of random modulation in the particle distribution, and to the variation of this modulation passage to passage, incoherent radiation is emitted with intensity fluctuating passage to passage.

It has been shown (Stupakov, Zolotorev SLAC-PUB 7132, 1996) that by measuring the passage-to-passage variance of the radiation intensity in a part of the spectrum where the emission is incoherent, the bunch length can be measured.



A non-destructive, remarkably simple bunch length measurement scheme has been developed that can be used in both circular and linear accelerators, with any kind of radiating process, including those cases where the very short length of the bunches makes difficult the use of other techniques.

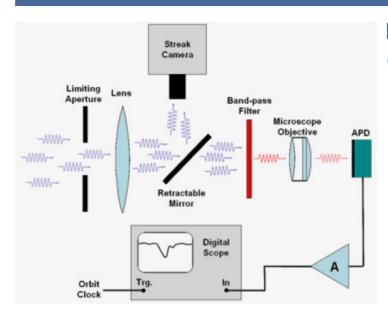


Highlights from the 13th BIW



Beam Absolute Bunch Length by Radiation Fluctuation Analysis





By using a bandpass filter with bandwidth σ_{ω} and measuring the passage to passage relative variance δ^2 of the radiation intensity, the bunch length σ_{τ} can be measured from:

$$\delta^2 = 1/\sqrt{1 + 4\sigma_{\tau}^2 \sigma_{\omega}^2}$$

The method has been successfully demonstrated at the Advanced Light Source using synchrotron radiation from a dipole.

Sannibale et al. PRST-AB 12, 032801 (2009).



Optical Diffraction Radiation Interferometry

DIFFRACTION RADIATION

- ➤ DR is produced by the interaction between the EM fields of the traveling charge and the conducting screen
- The extension of the electromagnetic field of a relativistic particle is a flat circle of diameter $\gamma \lambda/2\pi$,
- > The radiation intensity is

$$I \propto e^{-\frac{2\pi a}{\gamma \lambda}}$$

➤ DR impact parameter is

$$\frac{\gamma\lambda}{2\pi}$$
 \rightarrow if a

$$>> \frac{\gamma \lambda}{2\pi}$$
 No radiation
$$\cong \frac{\gamma \lambda}{2\pi}$$
 DR
$$<< \frac{\gamma \lambda}{2\pi}$$
 TR

Excellent candidate to measure beam parameters **parasitically**

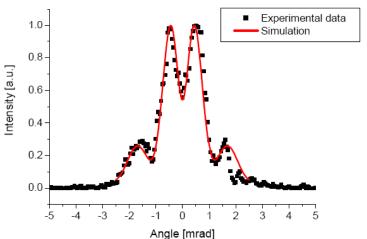


Optical Diffraction Radiation Interferometry

ODR ANGULAR DISTRIBUTION MEASUREMENT

The background has been a severe limitation (even at KEK) for a detailed and quantitative reconstruction of the beam parameters from the ODR angular distribution. A big effort has been done in order to subtract the SR background via software.

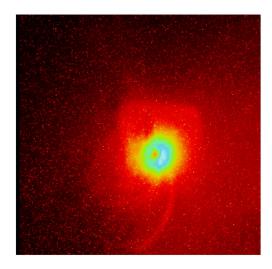
This allowed us to prove a good qualitative agreement between the experimental data and the simulations.



E. Chiadroni et al.,

Non-intercepting electron beam transverse diagnostics
with optical diffraction radiation at the DESY FLASH facility

NIM B 266 (2008) 3789–3796



Beam transport optimization

- > 0.7 nC
- > 25 bunches
- > 2 s exposure time
- $> E_{beam}$ (nominal) = 680 MeV
- > 800 nm filter and polarizer in

Simulation parameters:

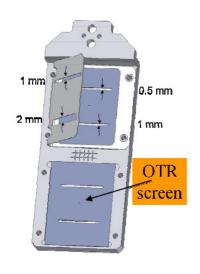
- > a = 0.5 mm
- > Gaussian distributed beam
- $\sigma_v = 80 \, \mu m$
- $> \sigma_{v}^{2} = 125 \,\mu rad$
- $E_{\text{beam}} = 610 \text{ MeV}$



Optical Diffraction Radiation Interferometry

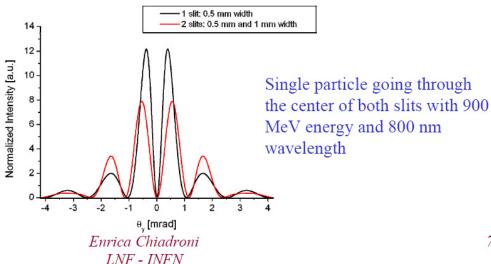
OPTICAL DIFFRACTION RADIATION INTERFEROMENTRY (ODRI)

To reduce the synchrotron radiation background, we mounted a stainless steel shield in front of our ODR screen, with a larger cut in it.



In the case of a wavelength of 800 nm and 1 GeV beam energy the 1 mm cut is not large enough to prevent the production of ODR in the forward direction, reflected by the screen and interfering with the backward ODR produced by the screen itself.

An ODR analogous of the Wartski interferometer used for OTR, with the difference that in this case the two interfering amplitudes are different in intensity and angular distribution





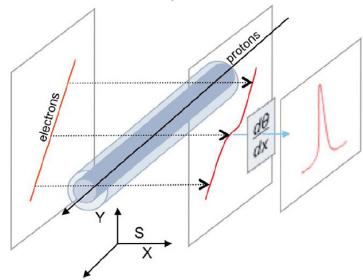
Principle

Look at the deflected projection of a tilted sheet of electrons due to the proton beam charge [1,2,3]

- Neglect magnetic field (small displacement of projection)
- Assume path of electrons is straight (they are almost straight)
- Assume net electron energy change is zero (if symmetric).

$$\Rightarrow \frac{d\theta_0(x)}{dx} = \int_L \frac{e}{mv^2} \frac{\delta(x,y)}{\varepsilon_0} dy \quad \text{or, take the derivative to get the profile}^{[3]}$$

Imperfections estimated at 5-10%, .



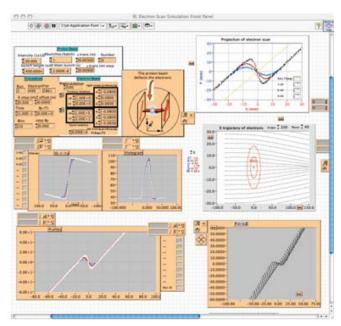


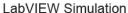


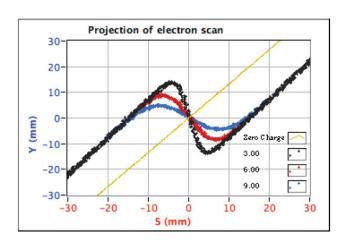
Principle

Approach:

- 1) Build an electron gun (75kV)
- 2) Shoot a tilted line of electrons through the proton beam
- 3) Use a fluorescent screen and video camera to get the projected curve





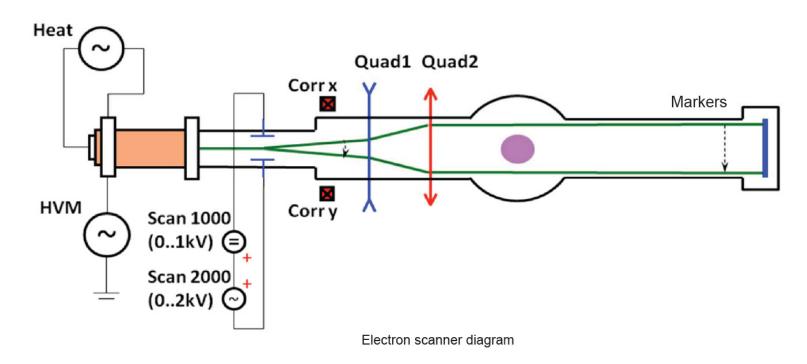


Projection of electrons (LabVIEW simulation)





Hardware: Electron Scanner



Electron Scanner made by Budker Institute of Nuclear Physics: Dmitriy Malyutin, Sasha Starostenko, Sasha Tsyganov Vacuum vessel by SNS.





Data

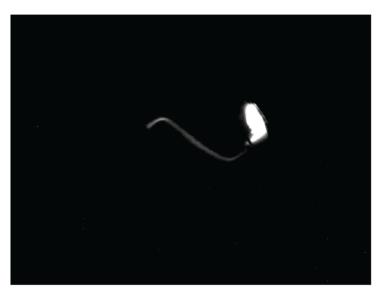


Image of horizontal curve

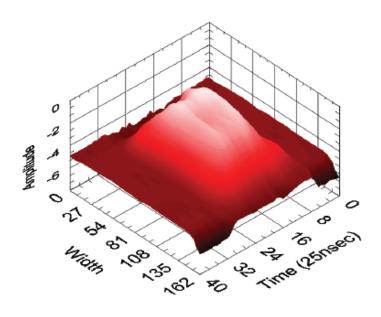


Image of vertical curve

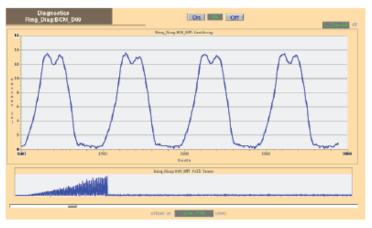




Data



3D plot of Turn 343.

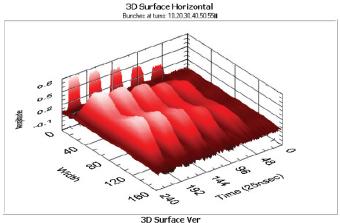


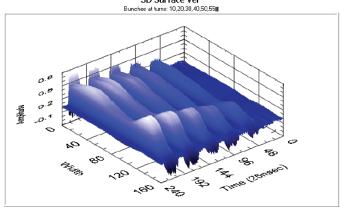
Ring BCM around Turn 343.



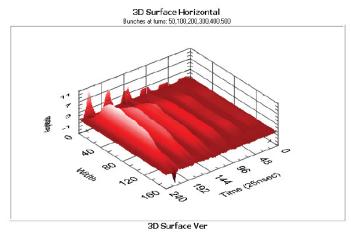


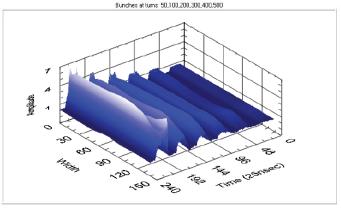
Data











3D plot of turns 50,100,200,300,400, and 500.





Conclusions

- Can deliver a profile for a 20 ns slice of beam anywhere in accumulation and anywhere along the bunch width.
- Can be operated completely parasitically during production.
- Profiles delivered look physical and quality (beam tilt, pulse to pulse jitter) is improving rapidly.
- The profiles the ELS is delivering look reasonable and physical, and the device has uniquely attractive features compared to traditional harp and wire scans.
- The profile shape are very sensitive to precise setting of the device hardware (quads and correctors).
- In it's current state, aperture is too small for large beams, and profiles not available for very low (< 2uC) or very high (> 9uC) intensities.

This is a great success for a new device.

Once minor issues are addressed, there will be a strong case for upgrading to a full tomography device.

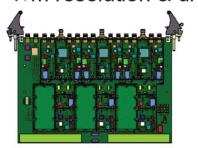




Cavity & Button BPM Electronics (PSI Designs)

Undulator RFFE

- 3.3GHz (cavity BPM)
- IQ demodulation
- Requirements: Sub m resolution & drift



ADC Mezzanine

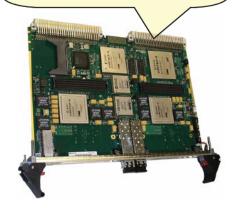
- Six 16-bit ADCs
- 160Msps



FPGA Carrier Board

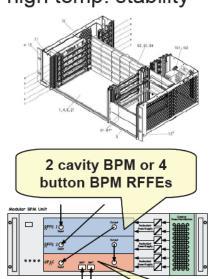
- Virtex-5 FPGAs
- Flexible interfaces:
 1-5Gbit Rocket IO,
 VME, VXS, Ethernet
- Two mezzanines:
 500-pin connectors

Low-cost version of IBFB carrier board (no DSPs, ...), used for all E-XFEL BPMs



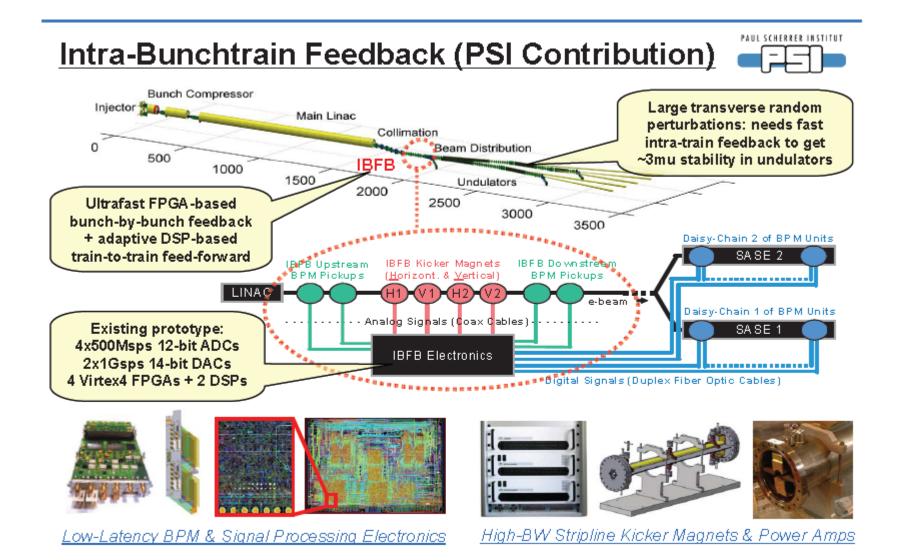
Modular BPM Unit

Crate: customized power, backplane & cooling: low noise, high temp. stability



FPGA/ADC board









BPMs for E-XFEL

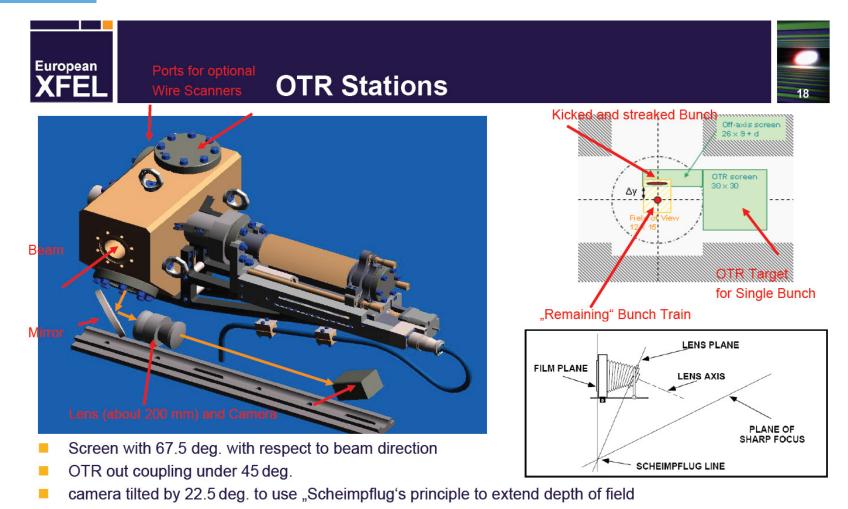


BPM Type	Number	Diameter	Single Bunch Resolution
Standard Button BPM	228	40 mm	50 µm
"cold" BPM (Button, 30% Re-entrant Cavity)	104	78 mm	50 µm
Precision BPM (Cavity)	117	10 mm	1 µm
Precision BPM (Cavity)	12	40 mm	1 µm

- Collaboration between PSI, CEA and DESY
 - PSI will provide Electronics, except RF front-end of the Reentrant Cavity BPM
 - DESY will provide Mechanics, except Reentrant Cavity BPM
 - CEA will provide Reentrant Cavity BPM, incl. RF front-end



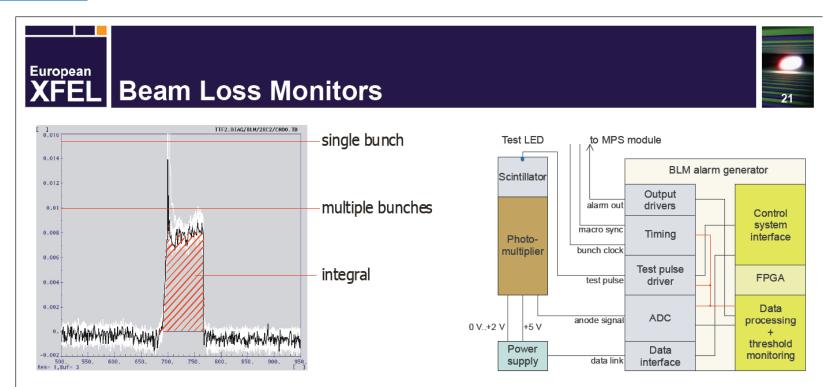




- 1: 1 reproduction scale
- Resolution requirement 10 30 μm (depending on section)
- Wire scanner ports for optional wire scanners
- Prototype test at FLASH scheduled
- Collaboration with IHEP, Protvino







- Good experience with FLASH system
- Stay with plastic scintillators and photomultipliers
- about 250 channels distributed over the entire machine
- Reengineering of the FLASH system
 - go to a pure digital readout and FPGA processing
 - AMC, µTCA form factor
- Collaboration with IHEP, Protvino

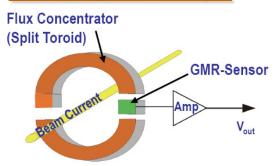




Beam Diagnostic Developments for FAIR

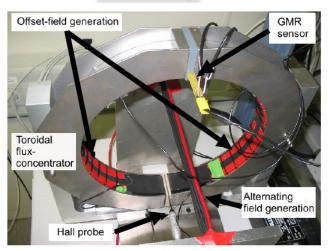
Beam Current Measurement - Novel DCCT

Measurement Principle

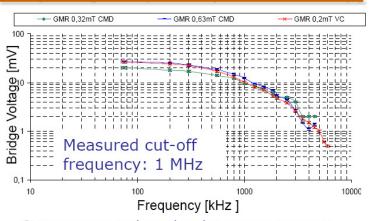


- idea: clip-on Amperemeter design
- split toroid to allow dismounting before bake-out
- soft-magnetic flux concentrator (amorph. VITROVAC®)
- gap with induction of 80 μT @ 1 A beam current
- sensitive GMR (Giant Magneto Resistance) magnetic field sensor (resolution: $10^{-9} \text{ T/}\sqrt{\text{Hz}}$) \rightarrow used for harddisks

Test Setup



Frequency Response (GMR+Pre-amp)



 Reasons: induced voltages in sensor material, eddy currents in NiFe-layer









Beam Diagnostic Developments for FAIR

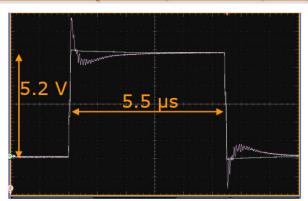
2nd Iteration: Combined ACT-DCCT-System

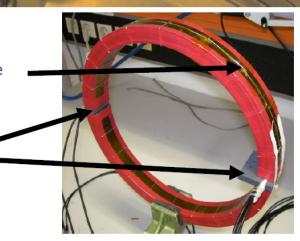
NDCCT Version 2.0

- add single ACT-winding to core
- add high-pass filtered ACT-branch to electronics
- use 2 GMR sensors (1 sensor per gap)

New: dc and ac measurement on single core!

Pulse Response (Combined System)





<u>Improvements</u>

- For S/N=2: 88 nV/√Hz
- calculated resolution threshold: 220 μΑ
- Bandwidth: $f_{co} \approx 13 MHz$

Ongoing R&D

- design & layout of readout electronics
- investigation of NDCCT rf- characteristics





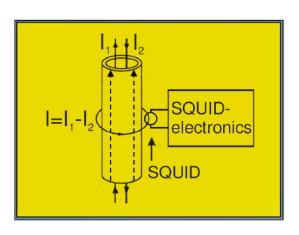


Intensity and Profile Measurements [...] in CSR



Cryogenic Current Comparator (CCC) Principle



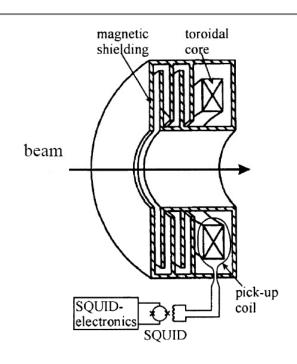


CCC (Harvey 1972):

- Uses Meissner-effect and SQUID for I₁/I₂ measurement
- If $I_1 \neq I_2$ magn. field produces compensation current
- Magnetic flux through SQUID → voltage change

For charged particle beams:

$$I_{comp} = I_1 - I_2 = I_{beam} - 0$$
 (position independent)



- SC shielding for non-azimuthal fields
- SC pickup coil with toroidal core ($\mu_r \approx 50000$)
- Low noise, high performance DC SQUID control electronics (FSU Jena)



Intensity and Profile Measurements [...] in CSR



Optimisation of CCC Performance



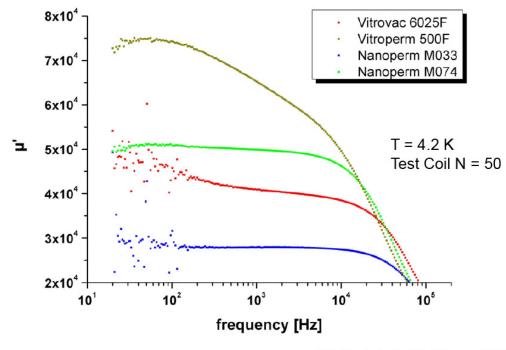
Achievements so far: $250 pA/\sqrt{Hz}$, BW = 0...50 kHz at GSI (A. Peters et al. 1999) \Rightarrow TARN II $40 pA/\sqrt{Hz}$, BW = 0...70 kHz at test setup for DESY (W. Vodel et al. 2007)

Limitations of the system:

- Mechanical vibrations
- Magnetic shielding
- Noise from toroidal core
- SQUID intrinsic flux noise
- Electronics (amplifier input noise, crosstalk etc.)
- Slew rate / core mat. (BW)

Current detection limit from pickup coil:

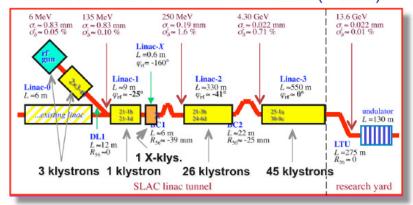
$$I_{S} = \frac{2\pi\sqrt{k_{b}TL}}{\mu_{0}\mu_{r}f(R_{a}, R_{i}, b)}$$



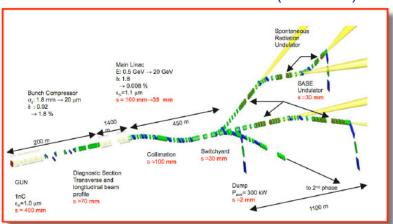
W. Vodel, R. Geithner (FSU)

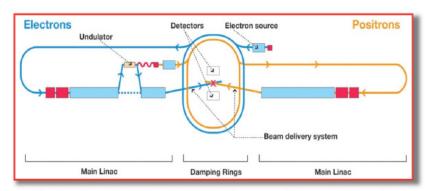


LCLS XFEL: 13.6GeV NC linac (~1km)

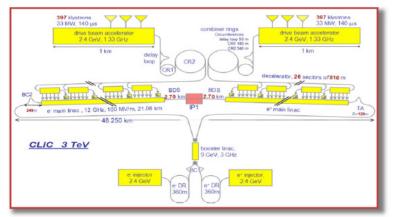


EU-XFEL: 20GeV SC linac (~1.6km)





ILC: 2x 250GeV SC linacs (2x ~11km)

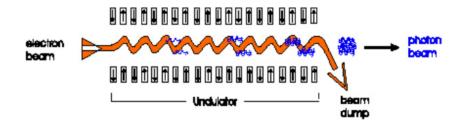


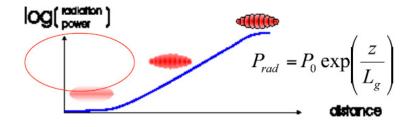
CLIC: 2x 1.5TeV NC linacs (2x ~21km) (+ drive beam linacs)



XFEL Principle

Linac based Self Amplification of Spontaneous Emission (SASE) Free Electron Lasers (FELs) in the X-Ray regime (~0.85 - 60 Å)





Gain Length
$$L_G = \frac{1}{\sqrt{3}} \left(\frac{I_A \gamma^3 \sigma_r^2 \lambda_u}{4\pi \hat{I} K^2} \right)^{\frac{1}{3}}$$

Electron bunch modulated with its own synchrotron radiation field

- ⇒ micro-bunching
- ⇒ more and more electrons radiate in phase until saturation is reached

Requires excellent electron beam quality:

- low (slice) emittance
- · low energy spread
- extremely high charge density

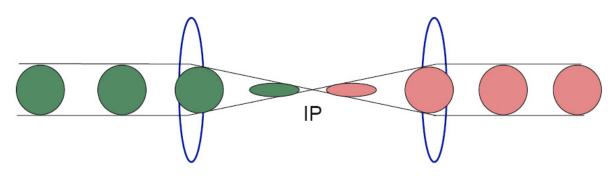
Winni Decking





Colliders

Squeeze as much as possible to get the highest luminosity



Luminosity (cm⁻² s⁻¹)
$$L \approx \frac{n_b N^2 f_{rep}}{A} H_D$$

 $n_b = \text{bunches / train}$

N = particles per bunch

 f_{rep} = repetition frequency

A = beam cross-section at IP

 H_D = e+/e- enhancement factor

Requires excellent beam quality:

- Extremely high particle densities
- Bunches colliding at the focal point



Linac stability expectations

FEL Photon Source

- Photon beam pointing accuracy and stability
- Photon wavelength, intensity (brilliance)
- Photon bunch arrive time relative to a fiducial (pump probe experiments)

Linear Collider

- Energy precision and stability
- Luminosity
- Integrated luminosity

Particle density

Relative pointing accuracy at the IP

Relative arrival time at the IP

Small beams (nm)
Low emittance (nm-rad)



Very tight tolerances
High-resolution diagnostics



Instrumentation Req. for Dif. Accelerator Types

III. Light Sources

The Standards: beam energy, intensity, tunes etc ... not really a problem

BUT measurement & control of beam size & beam orbits is a real challenge

photon flux
$$F = \frac{number\ of\ photons}{s*0.1\%BW*A}$$

brilliance
$$B = \frac{F}{4\pi^2 \varepsilon_x \varepsilon_y}$$

emittanz ≈ 1 nm rad (PETRA 3)

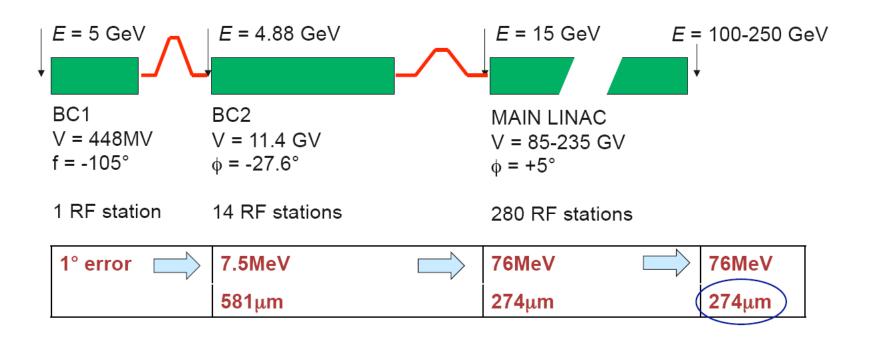
 $\varepsilon_v \approx 1\% \ \varepsilon_x$

ϵ_x [nmrad]	E [GeV]	ϵ_x/E^2
1.1	(A) (A) (A) (A)	
4.4	2.4	0.763
A 7	2.4	1.215
6	1.9	1.66
18	3	2
9	1.5	4
41	2.5	6.56
I 450	4.5	22.2
	A 7 6 18 9 41	6 1.9 18 3 9 1.5 41 2.5

Parameters of some synchrotron light sources



Sensitivity to a 1 degree error in BC1 Phase (1 klystron)



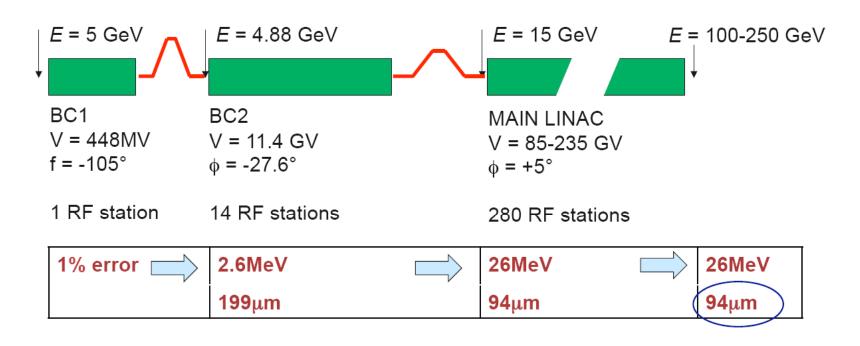
For bunch to arrive within 70μm of IP: BC1 phase error must be less than 0.25 deg

N. Walker





Sensitivity to a 1% error in BC1 Amplitude (1 klystron)



For bunch to arrive within 70μm of IP: BC1 amplitude error must be less than 0.74%

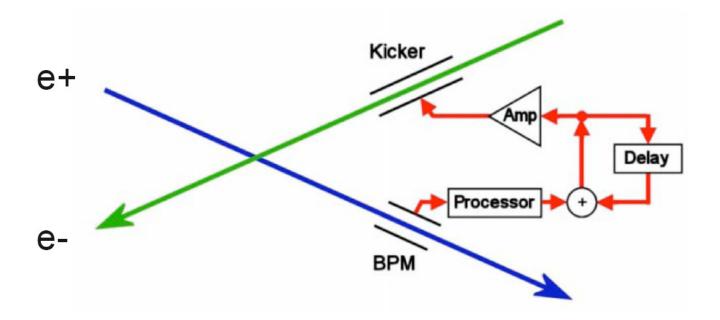
N. Walker



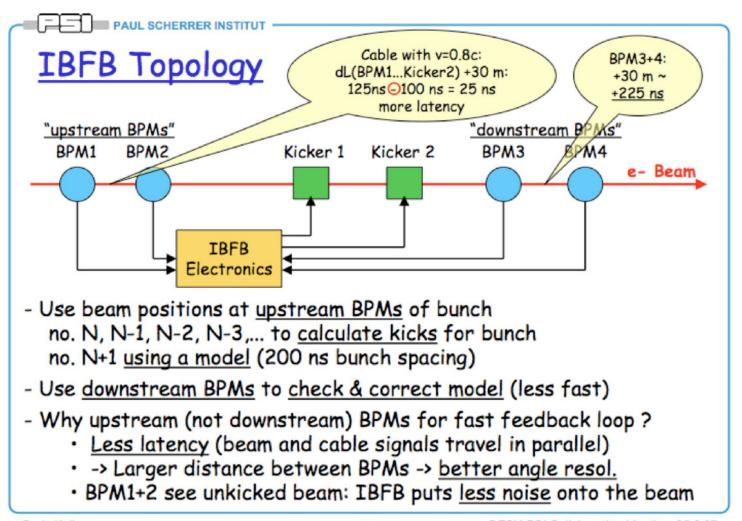


IP beam-based feedback for LC

- Maximize luminosity by correcting trajectory within bunch train
- Kicker and bpm must be close together for fast correction rate
- Need very high precision single-pass bpms







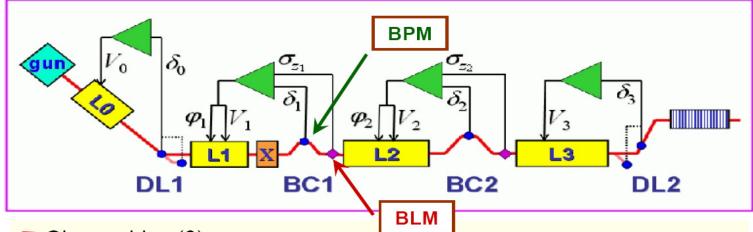




Stanford Linear Accelerator Center

Stanford Synchrotron Radiation Laboratory

LCLS longitudinal feedback system schematic



- Observables (6):
 - Energy: E_0 (at DL1), E_1 (at BC1), E_2 (at BC2), E_3 (at DL2)
 - Coherent Radiation power \rightarrow bunch length: $\sigma_{z,1}$ (at BC1), $\sigma_{z,2}$ (at BC2)
- Controllables (6):
 - Voltage: V₀ (in L0), V₁ (in L1), V₂ (effectively, in L2)
 - Phase: $\mathbf{φ}_1$ (in L1), $\mathbf{φ}_2$ (in L2), $\mathbf{φ}_3$ (in L3)

ICFA FLS, May 15-19, 2006 Hamburg, Germany Patrick Krejcik pkr@slac.stanford.edu







Poster highlights

MOPD11: INVESTIGATION OF PRECISE PIPELINE-TYPE ADCS IN A BURST REGIME FOR A SINGLE-SHOT BPM

MOPD13: IMPLEMENTATION OF AN FPGA-BASED LOCAL FAST ORBIT FEEDBACK AT THE DELTA STORAGE RING

MOPD42: PROFILEVIEW - A DATA ACQUISITION SYSTEM FOR BEAM INDUCED FLUORESCENCE MONITORS

TUPB02: BEAM INDUCED FLUORESCENCE MONITOR & IMAGING SPECTROGRAPHY OF DIFFERENT WORKING GASES

TUPB12: BEAM TEST OF THE FAIR IPM PROTOTYPE IN COSY

TUPB13: IR PHOTON ARRAY DETECTOR FOR BUNCH BY BUNCH TRANSVERSE BEAM DIAGNOSTICS

TUPB18: VIMOS, BEAM MONITORING FOR SINQ

TUPB24: BEAM HALO MONITOR USING DIAMOND DETECTOR FOR INTERLOCK SENSOR AT XFEL/SPring-8

TUPB35: VELOCITY OF SIGNAL DELAY CHANGES IN FIBRE OPTIC CABLES

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THANK YOU.