

# An overview on Beam Loss Monitoring

- The BLM business:
  - ▶ Definitions
  - ▶ Detector technologies
  - ▶ Machine protection
  - ▶ Location of BLMs
- Challenges on Beam Loss Monitoring



27-06-2014 BI seminar

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# Some references



- “Beam Loss Monitors”, K. Wittenburg. CERN Accelerator School 2008
- “Beam Loss Monitors”, L. Froelich . ERL instrumentation workshop 2008  
Accelerator School 2014
- “Beam instrumentation and diagnostics”, P. Forck. Joint University 2014
- “Overview of BLM technology”. K. Wittenburg. 3rd oPAC Topical Workshop  
on Beam Diagnostics 2014
- “Optical fiber based loss Monitors for electron storage rings”. T. Obina.  
IBIC2013
  
- “Beam instrumentation and diagnostics. P. Strehl
  
- Particle Data Group (<http://pdg.lbl.gov/>)
  
- The Beam Loss/Beam diagnostics sessions of: IBICXX, iPACXX, HBXX, ...



# The BLM business

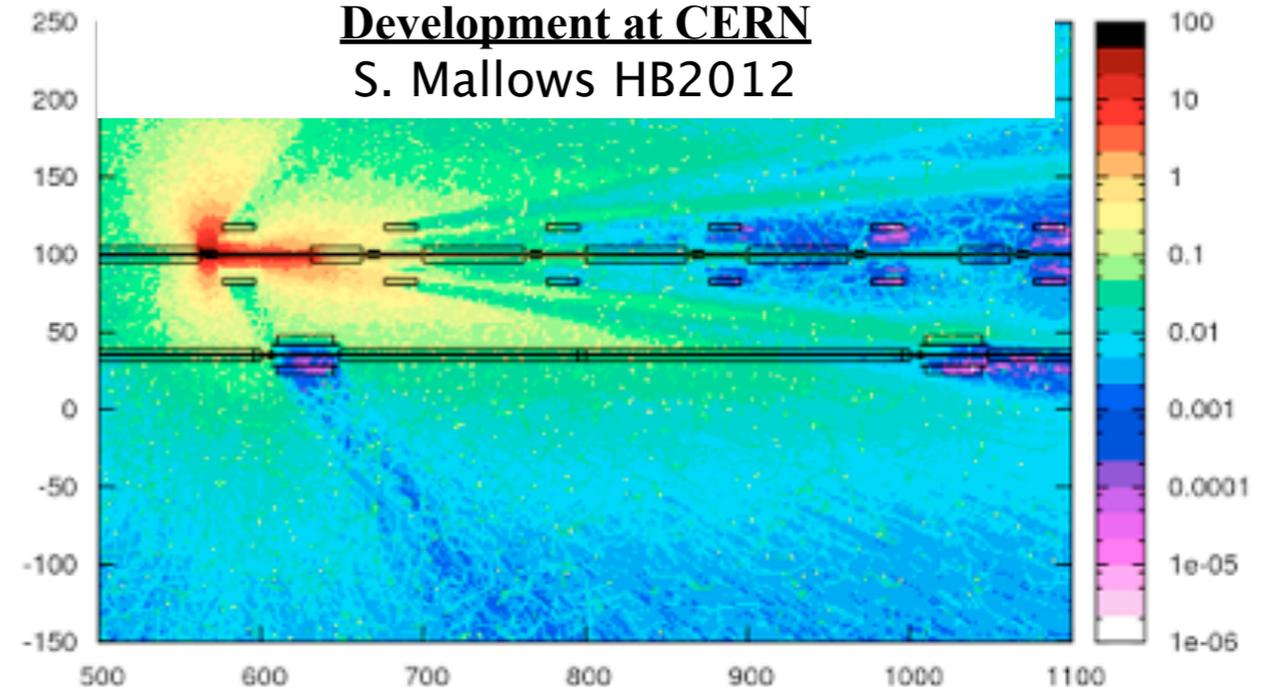
# Definitions



- **Beam Losses:**

- Particles deviating from the design orbit may eventually hit the aperture limit and be “loss” from the beam
- The impact of particles in the vacuum chamber produces particle showers

**Fiber Based BLM System Research and Development at CERN**  
S. Mallows HB2012



- **Beam Loss Monitoring system**

- Ionizing radiation detectors located around the accelerator
- What kind of detectors? Where do I put them? How do I extract/process the signals?

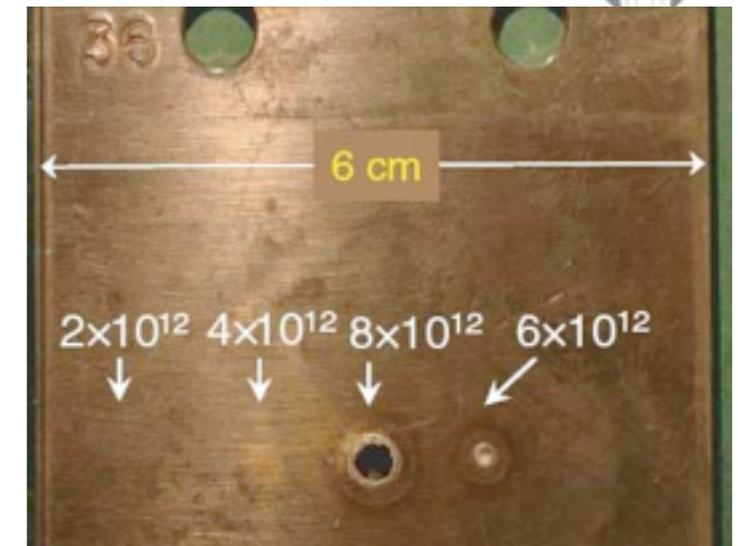


# Main goals



- Machine protection

- The impact of a high power beam into a single location may have catastrophic consequences on equipment.
- Aborts (safely extracts from ring or inhibits further injections in a LINAC) the beam when the BLM signals exceed a predefined threshold
- Thresholds established based on: Damage Level (Typically to vacuum chamber.beam pipe) or Quench levels in superconducting magnets

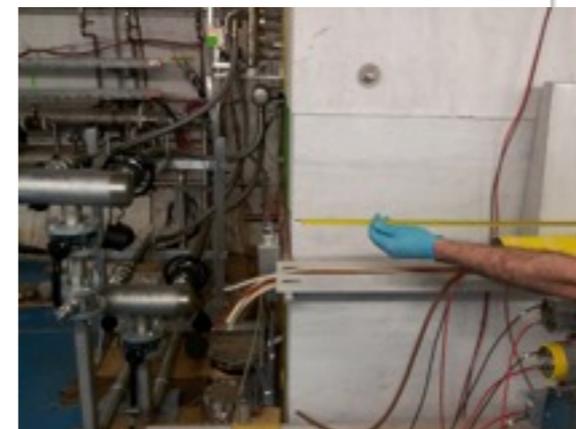
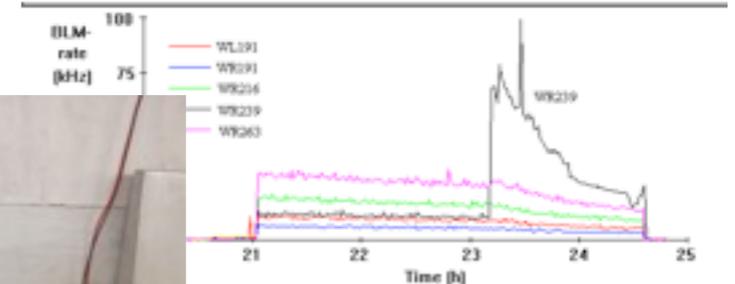
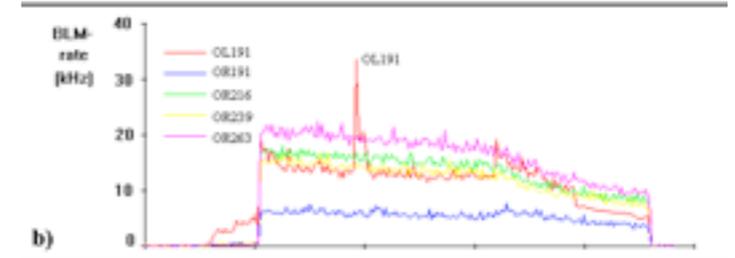
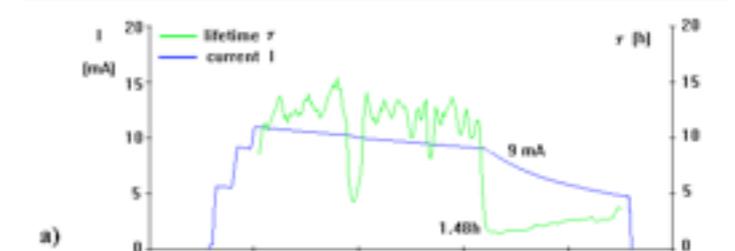


- Beam diagnostics

- Beam tuning
- Halo measurement

- Keep activation levels low

- Production of “radioactive waste”
- Protection against human hazard



- Beam Loss categorization

- Irregular beam losses

- ▶ They are avoidable but sometimes tolerated
- ▶ RF trips
- ▶ vacuum leaks
- ▶ Injection losses
- ▶ Obstacles in the beam

- Regular beam losses: Normally there are controlled although not avoidable

- ▶ Debris from interaction point (collider)
- ▶ Intentionally produced losses (for beam setup)
- ▶ Losses at aperture imitations (collimation)

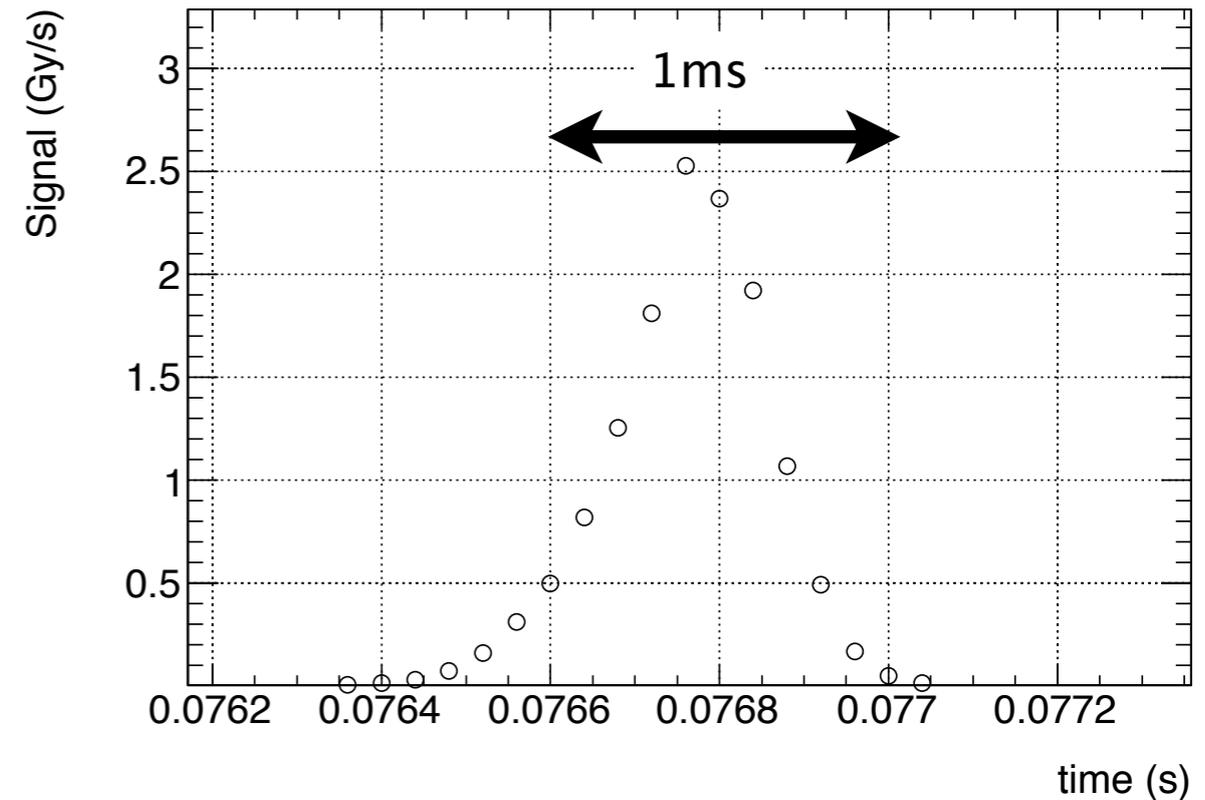
# Type of losses



- Beam Loss categorization

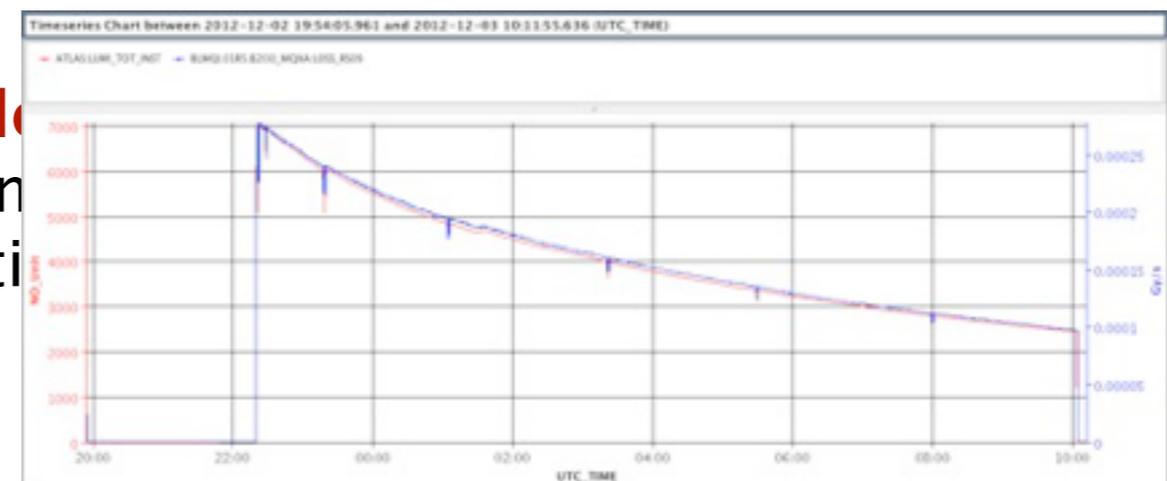
- Irregular beam losses

- ▶ They are avoidable but sometime
- ▶ RF trips
- ▶ vacuum leaks
- ▶ Injection losses
- ▶ **Obstacles in the beam**



- Regular beam losses: Normally there are controlled although not avoidable

- ▶ **Debris from interaction point (collider)**
- ▶ Intentionally produced losses (for beam)
- ▶ Losses at aperture imitations (collimati



# Source of BLM signal



- Ionization

- Energy loss by ionization described by the Bethe-Bloch formula
- Concept of Minimum Ionizing Particle

$$dE/dx_{MIP} = (1-5) \text{ MeV cm}^2 \text{ g}^{-1}$$

- Scintillation

$$Y = dL/dx = R dE/dx$$

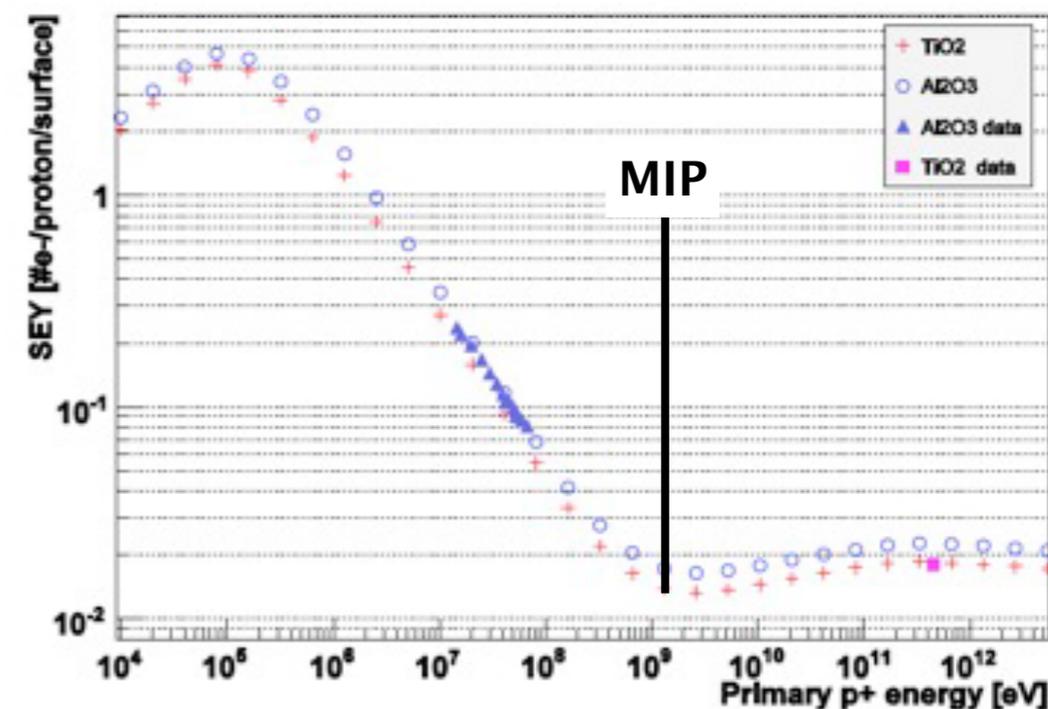
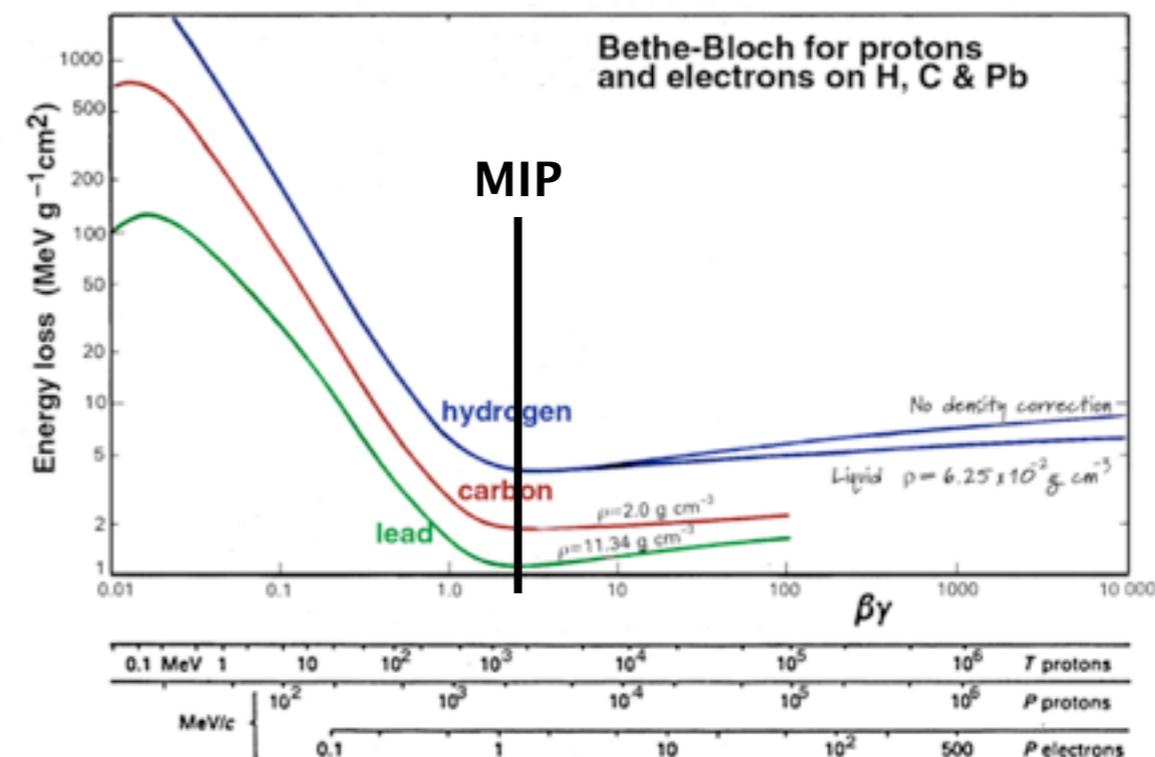
- Secondary emission

$$Y_{MIP} = (0.01-0.05) \text{ e/primary}$$

- Cherenkov light

$$\text{photon yield} : \frac{dN}{dx} = 2 \cdot \pi \cdot \alpha \cdot \sin^2 \Theta \cdot \left( \frac{1}{\lambda_1} - \frac{1}{\lambda_2} \right)$$

$$\cos \Theta = \frac{1}{\beta \cdot n} \text{ with } \beta > 1/n; \alpha = 1/137.036 \text{ and } \lambda_{1,2} = \text{wavelength interval}$$



# Some useful units and conversion numbers



You will need this for doing  
BLM R&D....  
... on a piece of paper

- With the assumption

$$dE/dx_{MIP} = 2 \text{ MeV cm}^2 \text{ g}^{-1}$$

a conversion between energy deposition  
(dose) and number of crossing minimum  
ionizing particles can be provided

## Radiation units



1 Ci =  $3.7 \times 10^{10}$  decays per second  
Marie and Pierre Curie



1 Bq = 1 decay per second  
Henri Becquerel

Activity – how frequently  
radioactive materials disintegrates.



1 R =  $2.58 \times 10^{-4}$  C/kg (radiation ionizing 1 kg  
of dry air to this charge)  
Wilhelm Röntgen  
Exposure (how much of material was ionized)

1 Sv = 1 J/kg,  
takes into consideration  
biological impact (quality factor)



Louis Harold Gray 1 Gy =  $1 \frac{\text{J}}{\text{kg}} = 1 \text{ m}^2 \cdot \text{s}^{-2}$

1 Rad = 0.01 Gy  
Dose (how much energy was deposited)

Rolf Maximilian Sievert  
Dose equivalent (how much harm  
the deposited energy caused  
to human)  
1 Rem = 0.01 Sv

21 Managed by UT-Battelle  
for the U.S. Department of Energy

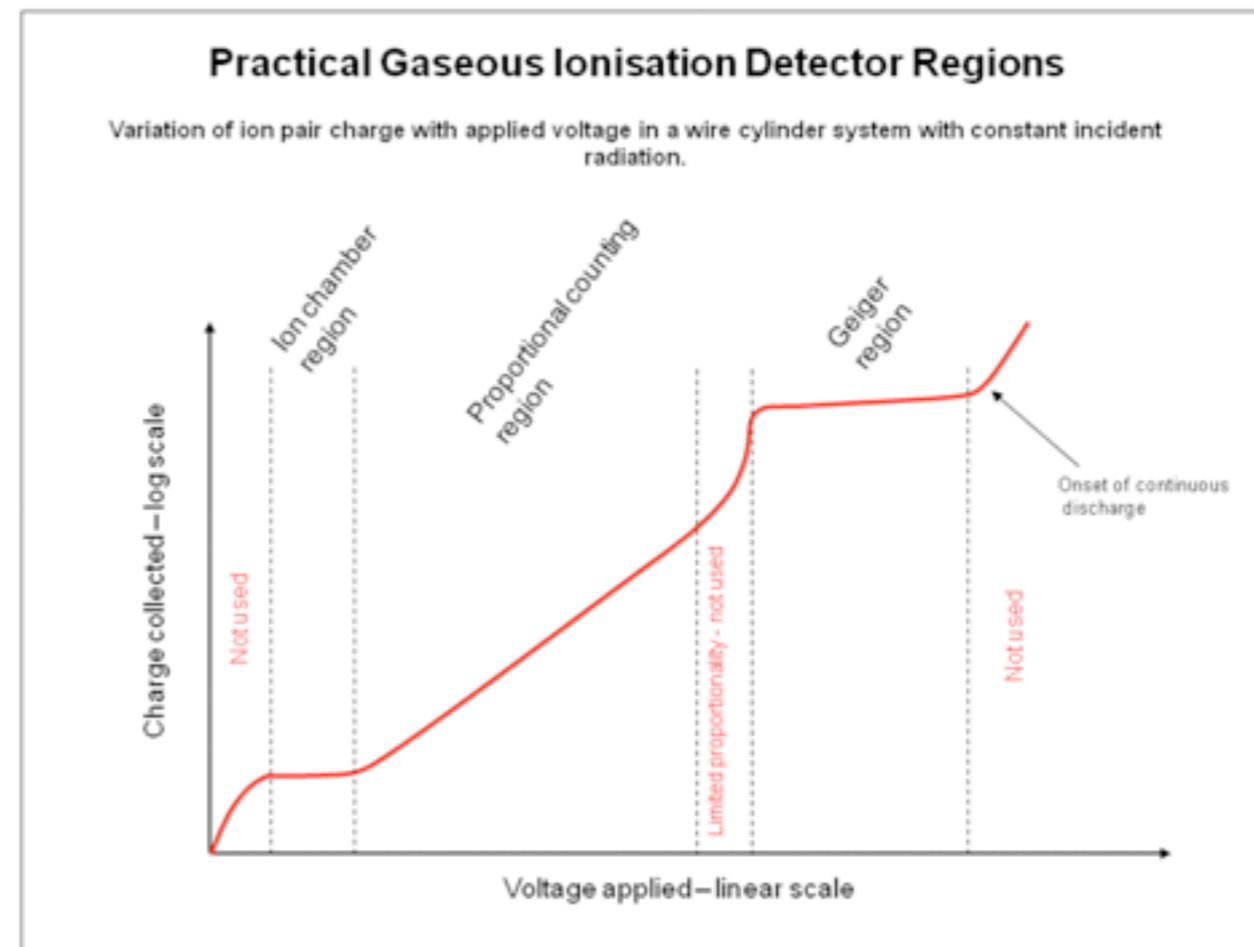
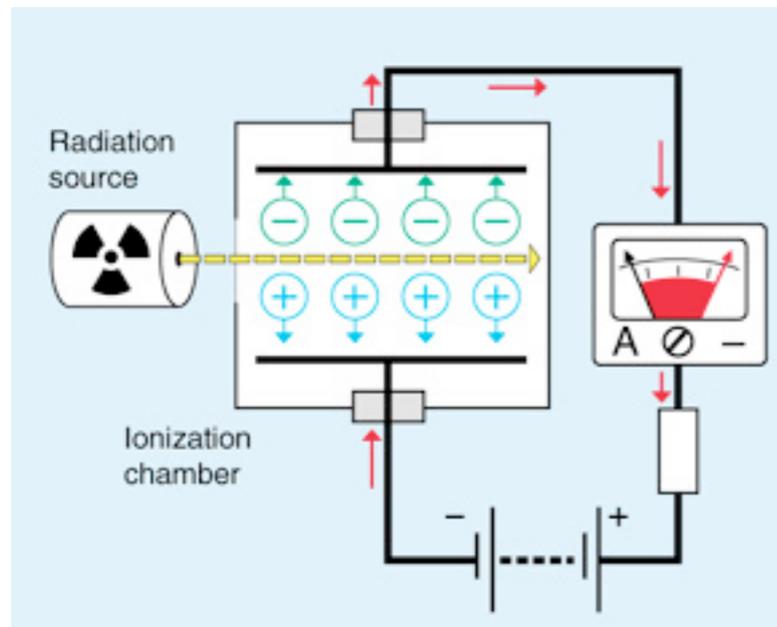
THINB01 BRW2010



$$1 \text{ Gy} = 1 \frac{\text{J}}{\text{kg}} \cdot \frac{\text{MIP} \cdot \text{kg}}{3.2 \cdot 10^{-10} \cdot \text{J cm}^2} = 3.1 \cdot 10^9 \frac{\text{MIP}}{\text{cm}^2}$$

## ● Ionization Chamber

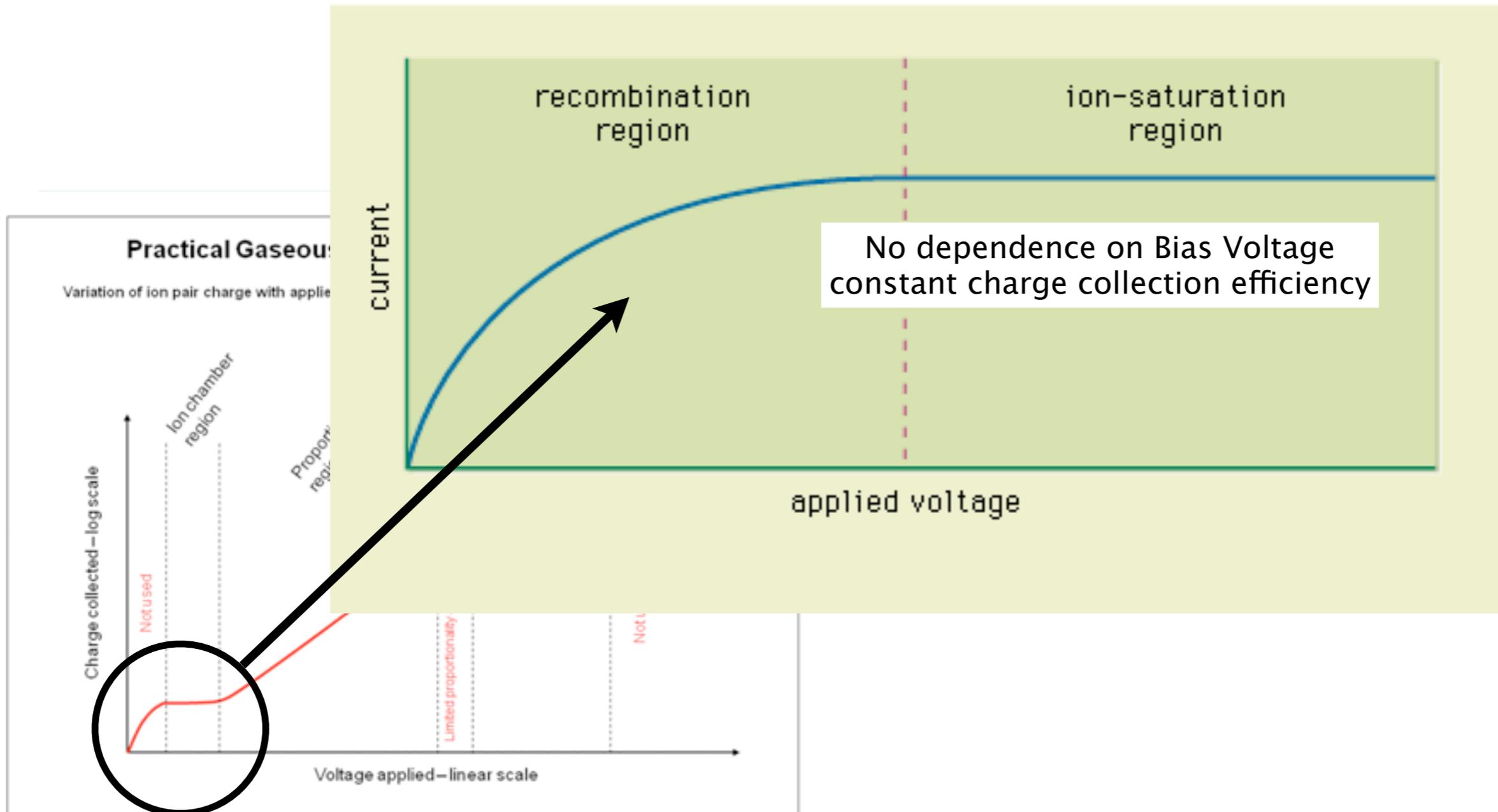
- Gas detector: Incident particle ionizes the active gas producing electrons (fast, a few  $\sim 100\text{ns}$ ) and ions (slow, a few  $100\mu\text{s}$ ) as charge carriers
- $(25-100)$  eV/pair depending on active gas



# Ionization based BLMs

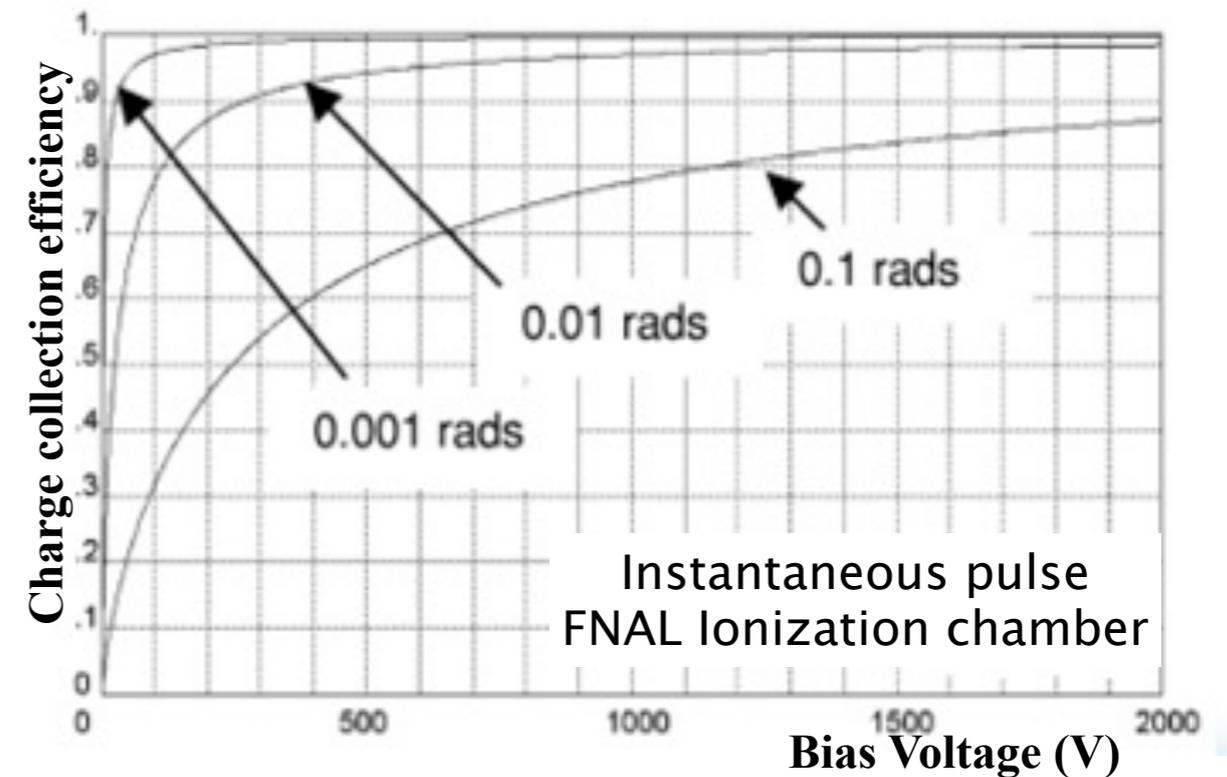


- Ionization Chamber



## • Ionization Chamber: Saturation

- Large number of e/ion pairs generated
- The field generated by ions shields the Bias field
- Charge recombination

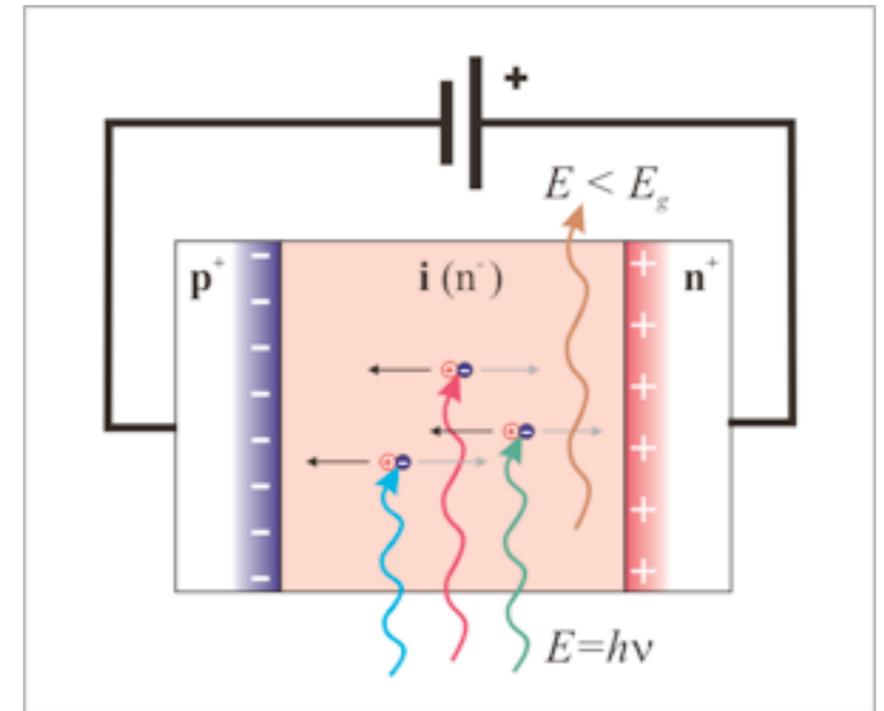


## PRO/CONTs:

- Very robust, radiation hard and require little maintenance
- No dependence on Bias Voltage
- Large dynamic range (up to  $10^{+8}$ - $10^{+9}$ )
- Slow (time resolution  $\sim 100\mu\text{s}$ )
- Saturation effects

- Semiconductor (PIN, diamond, Si,....). Solid state ionization chambers

- Incident particles produce electron/holes as charge carriers (3–10eV/pair)
- $t_{\text{hole}} \approx t_{\text{electron}} \sim 5\text{--}10\text{ ns}$
- smaller size



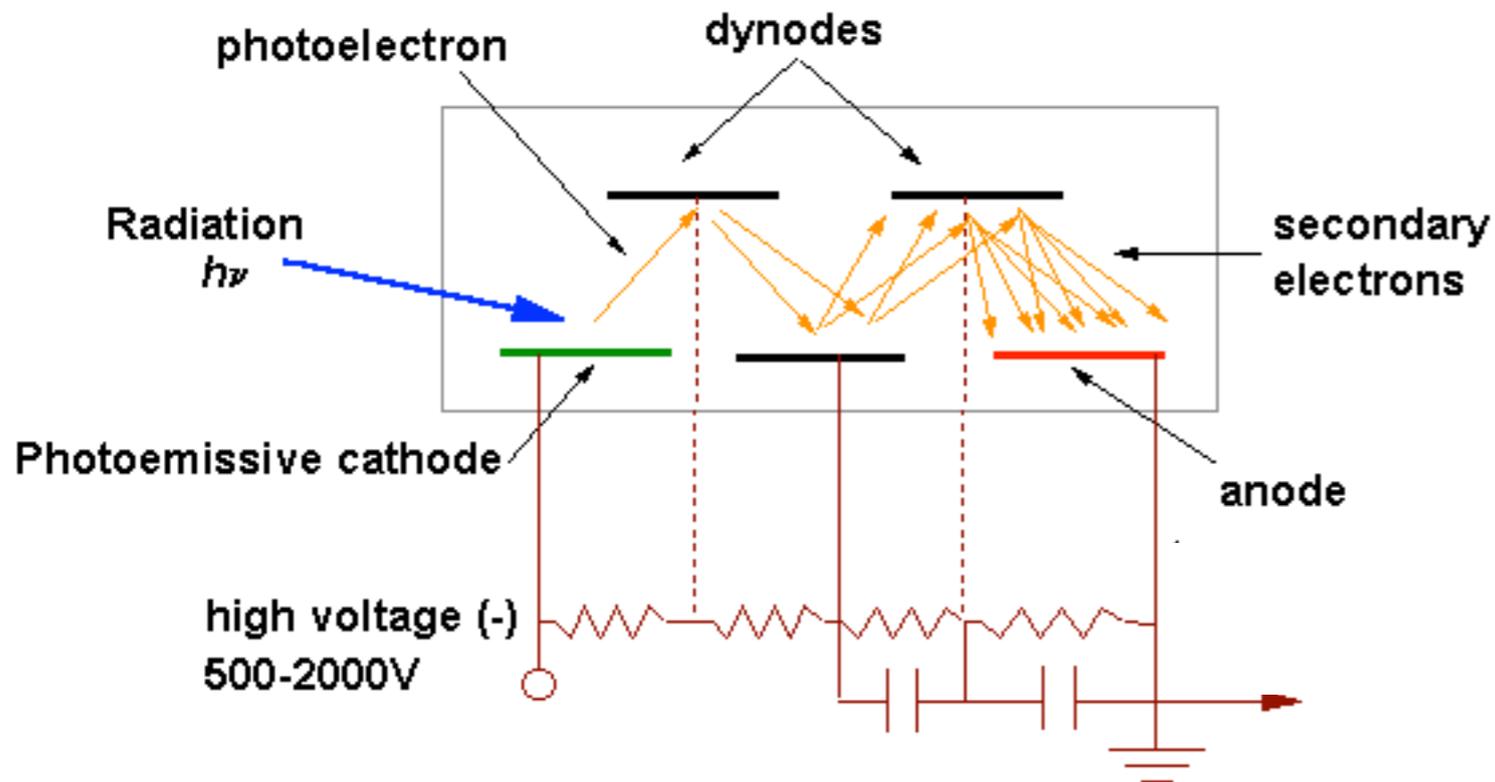
## PRO/CONTs:

- No dependence on Bias Voltage
- Fast (er) response (5–10 ns)
- Radiation hardness (1 MGy)

# Scintillation based BLMs



- Light produced by de-excitation of atomic/molecular levels
- Several types of scintillators
  - Inorganic crystals: NaI, CsI, ....
  - Organic (plastic): NE102, Anthracene,...
  - Liquid



- Light directed (via Waveguides) to photomultiplier tube

# Scintillation based BLMs



- What needs to be considered

- Photon Yield
- Collection Efficiency
- Photocathode quantum efficiency
- $G_{\text{PMT}} = \delta^n = (10^{+5} - 10^{+8})$  with:
  - ▶  $\delta$  (2–10) the number of secondary electrons
  - ▶  $n$  (8–15) the number of dynodes

Material	$R_s$ ( $\gamma/\text{MeV}$ )	$\rho$ ( $\text{g}/\text{cm}^3$ )
NaI	$8 \cdot 10^{+4}$	3.7
PbWO <sub>4</sub>	$2 \cdot 10^{+2}$	8.3
NE102	$2.5 \cdot 10^{+4}$	1.03
Antracene	$4 \cdot 10^{+4}$	1.025

## PRO/CONTs:

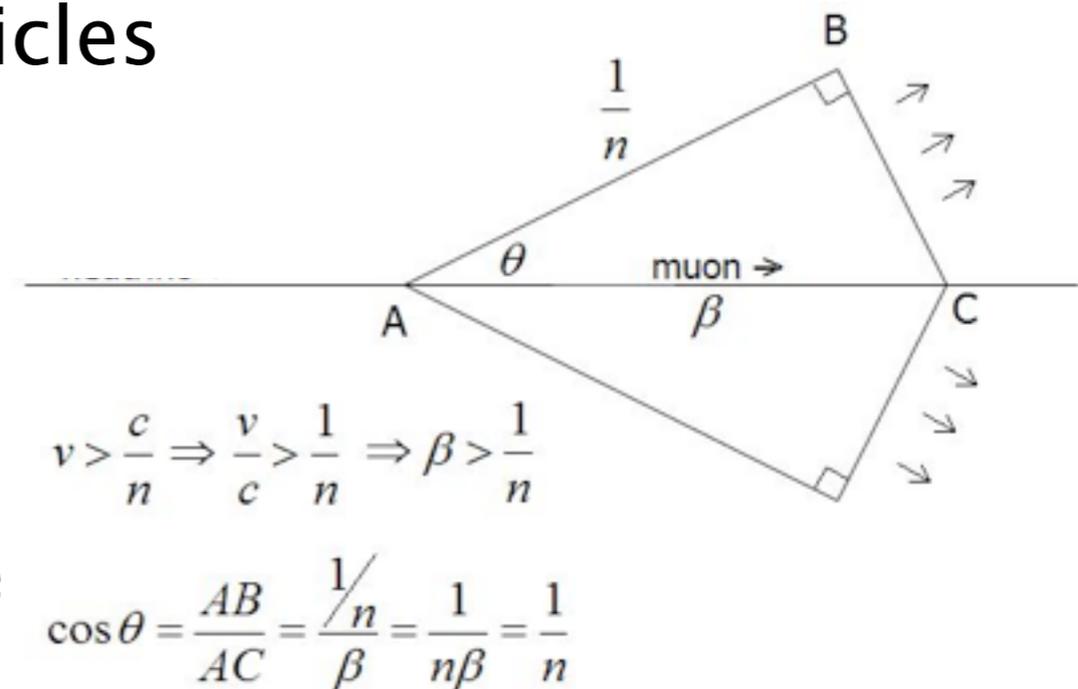
- High sensitivity
- Fast (5–10 ns) response (plastic/liquid)
- Slow (100ns–1us) response (inorganic)
- Limited radiation hardness (1MGy, 10MGy for liquids)
- Gain control of PMT

# Cherenkov light based BLMs



- Light produced by charged particles traveling at speed  $v > v_{\text{light}}$

- Photon Yield  $\sim 1000$  photons/mm (assuming quartz and a range 200–900nm)
- Cherenkov spectrum proportional to  $\lambda^{-2}$ . Only a small fraction on the light is usable



- As for scintillators we need to consider collection efficiency and quantum efficiency of photosensor

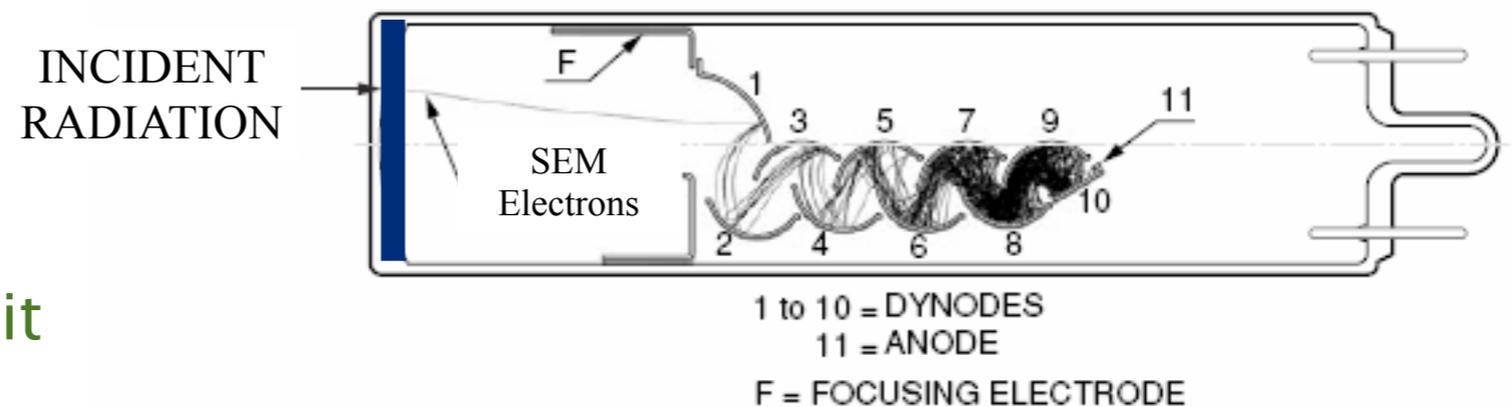
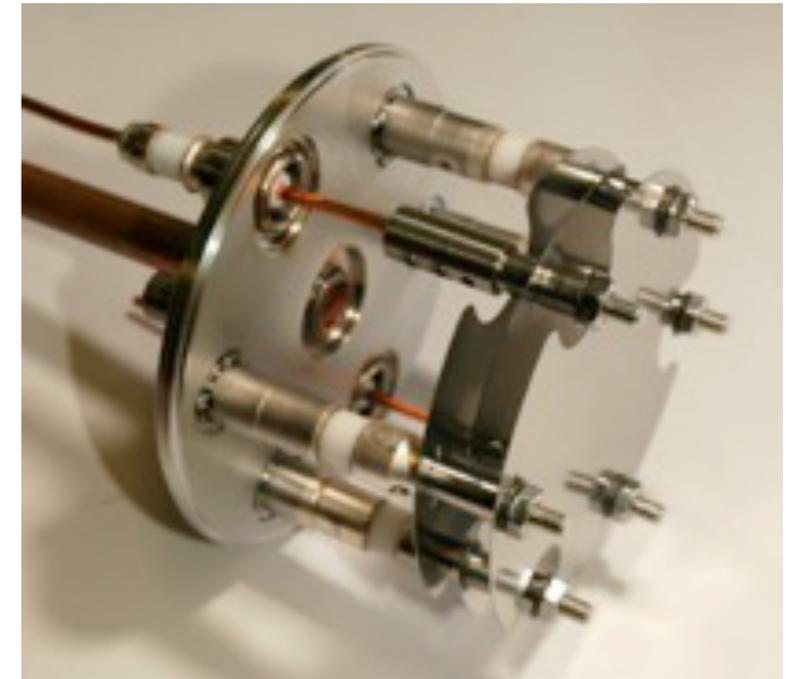
## PRO/CONTs:

- Insensitive to neutral radiation
- Fast (defined by your PMT or photosensor)
- Low sensitivity
- Limited radiation hardness (10–100MGy)

# Secondary emission based BLMs



- Sensitivity defined by the (1–5)% charges produced per primary
- Low possibilities:
  - With no amplification
    - ▶ Very low sensitivity
    - ▶ Needs current integration
  - With amplification
    - ▶ broadband current amplifiers
    - ▶ PMTs (sensitive to B fields)



## PRO/CONTs:

- Fast (<10ns). Electron transit
- Very linear
- Radiation hard (ideal for high radiation environments)
- Very low sensitivity
- Insensitive to neutral particles

# Sensitivity summary



- Ionization chamber:** **70  $\mu\text{C}/\text{Gy}$**   
 1 liter argon  
 $S \approx \text{active mass} \cdot \text{charge per ionization energy} \approx V \cdot \rho \cdot e / E_{\text{ion}} \approx 1 \text{ l} \cdot 1.8 \text{ g/l} \cdot e / 26 \text{ eV}$
  - Long ionization chamber:** **20  $\mu\text{C}/\text{Gy}$**   
 1 meter length, 1 cm radius, argon  
 $S \approx \text{active mass} \cdot \text{charge per ionization energy} \approx \pi r^2 \cdot L \cdot \rho \cdot e / E_{\text{ion}} \approx 314 \text{ cm}^3 \cdot 1.8 \text{ g/l} \cdot e / 26 \text{ eV}$
  - PIN diode:** **6  $\mu\text{C}/\text{Gy}$**   
 1  $\text{cm}^2$  surface, 100  $\mu\text{m}$  depletion depth  
 $S \approx \text{active mass} \cdot \text{charge per excitation energy} \approx A \cdot d \cdot \rho \cdot e / E_{\text{ion}} \approx 10 \text{ mm}^3 \cdot 2.3 \text{ g/cm}^3 \cdot e / 3.6 \text{ eV}$
  - Secondary emission monitor:** **500  $\mu\text{C}/\text{Gy}$**   
 100  $\text{cm}^2$  surface, 0.01 average secondary emission yield (SEY)  
 $S \approx \text{surface} \cdot \text{SEY} \cdot \text{electron charge} \cdot \text{density of primaries per dose} \approx A \cdot \text{SEY} \cdot e \cdot (\rho / (dE/dx))$   
 $\approx 100 \text{ cm}^2 \cdot 0.01 \cdot e \cdot 1 / (2 \text{ MeV} \cdot \text{cm}^2/\text{g})$
  - Aluminum cathode electron multiplier:** **5  $\mu\text{C}/\text{Gy}$**   
 10  $\text{cm}^2$  surface, 0.01 average secondary emission yield (SEY), tube gain  $10^5$   
 $S \approx \text{surface} \cdot \text{SEY} \cdot \text{electron charge} \cdot \text{density of primaries per dose} \cdot \text{gain} \approx A \cdot \text{SEY} \cdot e \cdot (\rho / (dE/dx)) \cdot G$   
 $\approx 10 \text{ cm}^2 \cdot 0.01 \cdot e \cdot 1 / (2 \text{ MeV} \cdot \text{cm}^2/\text{g}) \cdot 10^5$
  - PMT with organic scintillator:** **200  $\text{C}/\text{Gy}$**  ← **Radiation damage problematic!**  
 1 liter scintillator, 60% collection efficiency, 30% photocathode efficiency, tube gain  $10^5$   
 $S \approx \text{active mass} \cdot \text{photon yield per energy} \cdot \text{collection efficiency} \cdot \text{photocathode efficiency} \cdot \text{gain} \cdot \text{electron charge}$   
 $\approx V \cdot \rho \cdot Y \cdot C \cdot P \cdot G \cdot e \approx 1 \text{ l} \cdot 1 \text{ g/cm}^3 \cdot 1 / (100 \text{ eV}) \cdot 0.6 \cdot 0.3 \cdot 10^5 \cdot e$
  - Bare PMT (Čerenkov light):** **4  $\text{mC}/\text{Gy}$**   
 10  $\text{cm}^2$  surface, 1 mm thick, 30% photocathode efficiency, tube gain  $10^5$   
 $S \approx \text{active volume} \cdot \text{density of primaries per dose} \cdot \text{photon yield per length} \cdot \text{photocath. efficiency} \cdot \text{gain} \cdot \text{electron charge}$   
 $\approx A \cdot d \cdot \rho \cdot (\rho / (dE/dx)) \cdot Y \cdot P \cdot G \cdot e \approx 1 \text{ cm}^3 \cdot 1 / (2 \text{ MeV} \cdot \text{cm}^2/\text{g}) \cdot 260/\text{cm} \cdot 0.3 \cdot 10^5 \cdot e$
  - PMT with Čerenkov fiber:** **2  $\mu\text{C}/\text{Gy}$**   
 1 meter length, 100  $\mu\text{m}$  radius, 2% collection efficiency, 30% photocathode eff., tube gain  $10^5$   
 $S \approx \text{active volume} \cdot \text{density of primaries per dose} \cdot \text{photon yield per length} \cdot \text{coll. eff.} \cdot \text{photoc. eff.} \cdot \text{gain} \cdot \text{electron charge}$   
 $\approx \pi r^2 \cdot L \cdot \rho \cdot (\rho / (dE/dx)) \cdot Y \cdot C \cdot P \cdot G \cdot e \approx 31 \text{ mm}^3 \cdot 1 / (2 \text{ MeV} \cdot \text{cm}^2/\text{g}) \cdot 260/\text{cm} \cdot 0.02 \cdot 0.3 \cdot 10^5 \cdot e$
- Flexible gain → linearity and calibration problematic!

Detector material	Energy to create one electron [eV/e]	Number of [e / (cm MIP)] (depends on dE/dx and density)	Sensitivity S (for MIPs) [nC/rad]
Plastic scintillator	250–2500	$10^3 - 10^4$	$\approx 18 \cdot 10^3 \cdot (\text{PMT}_{\text{gain}})$ (1 ltr.)
Inorganic scint	50–250	$10^4 - 10^5$	$\approx 200 \cdot 10^3 \cdot (\text{PMT}_{\text{gain}})$ (1 L)
Gas ionization	22–95	$\approx 100$ (Ar, 1 atm., 20°C)	$\approx 600 \cdot (\text{Elec}_{\text{gain}})$ (1 L)
Semiconductor (Si)	3.6	$10^6$	$\approx 100 \cdot (\text{Elec}_{\text{gain}})$ (1 $\text{cm}^2$ PIN diode)
Secondary emission	2–5%/MIP (surface only)	0.02–0.05 e/MIP	$\approx 2 \cdot 10^3 \cdot (\text{PMT}_{\text{gain}})$ (8 $\text{cm}^2$ )
Čerenkov light	$10^5 - 10^6$	$\approx 10$ ( $\text{H}_2\text{O}$ )	$\approx 270 \cdot (\text{PMT}_{\text{gain}})$ (1 L)

# BLM choice and design considerations



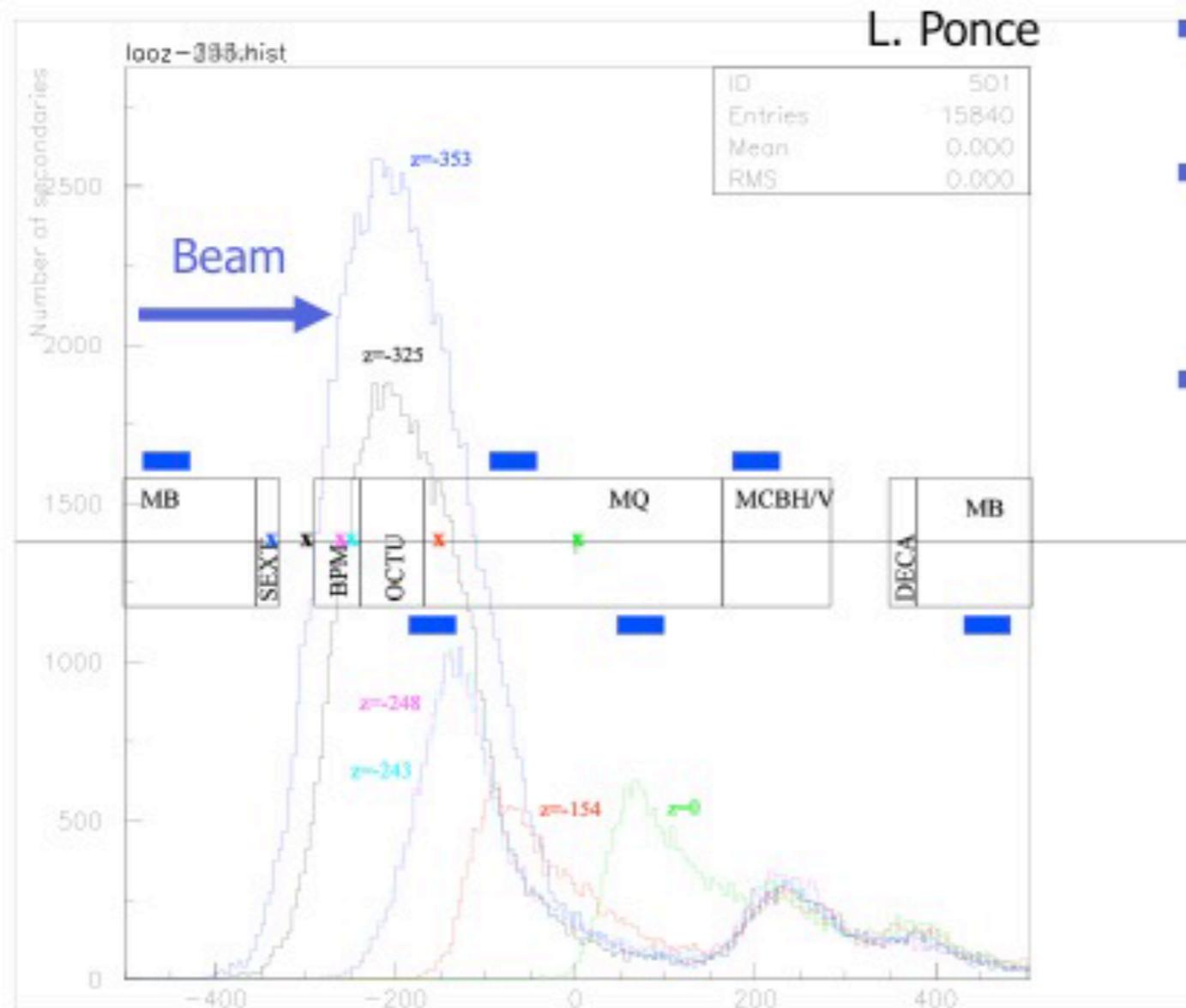
- No Universal rule!!!. Your BLM system will be designed for one particular machine
  - you CAN/SHOULD learn from BLM systems in previous machines but you won't be able to just copy any of them
  - Every accelerator is different: intensity, timing, radiation, length, normal conducting/superconducting ...
- To design a BLM system we need to consider:
  - Sensitivity
  - Dynamic range
  - Time response
  - Type of radiation
  - Shield-ability (from unwanted radiation)
  - Response to excessive radiation (saturation effects)
  - Physical size of BLMs
  - Test-ability
  - Calibration techniques
  - System end to end online test
  - Cost

# Location of BLMs



- Where is the ideal location of a BLM? Everywhere!!!!
- What's the best “realistic” location of a BLM?
  - Beam Physics. Particles more likely to be lost at:
    - ▶ Beam envelope maxima
    - ▶ Aperture limitations
    - ▶ Both
  - Comprehensive study of possible failure cases and consequences on beam
    - ▶ Miss-alignments
    - ▶ Asynchronous firing of kicker magnets
    - ▶ Some surprises may come once your machine is running
  - Simulations
    - ▶ Tracking codes (MADX, SixTrack,.....)
    - ▶ Monte Carlo simulation codes (Geant4, FLUKA, EGS5, ...)
    - ▶ Benchmarking experiments (i.e. don't trust your simulations)

## Particle Shower in the Cryostat



- Impact position varied along the MQ

- Black impact position corresponds to peak proton impact location

- Position of detectors optimized

- to catch losses:

- Transition between MB – MQ
- Middle of MQ
- Transition between MQ – MB

- to minimize uncertainty of ratio of energy deposition in coil and detector

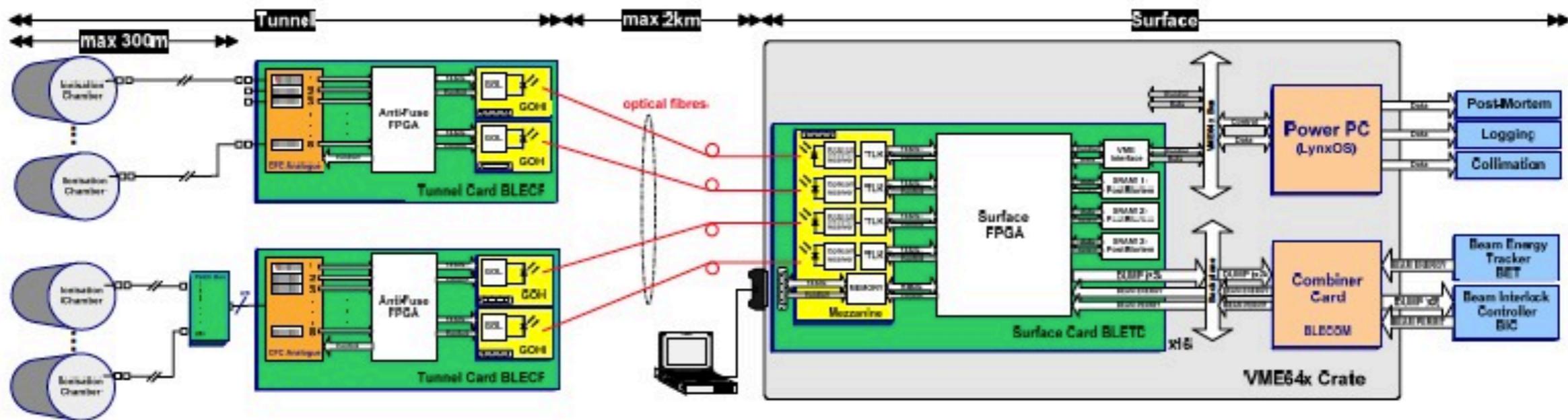
- Beam I – II discrimination

- Good probability that losses are seen by two BLM detectors

# Machine Protection: the LHC case



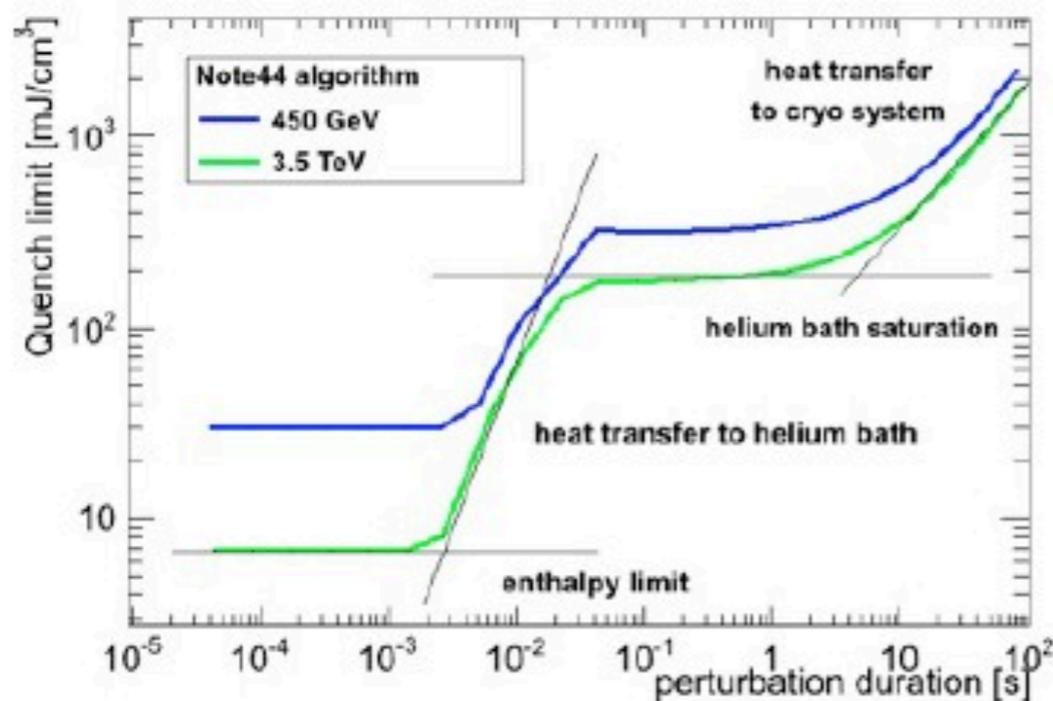
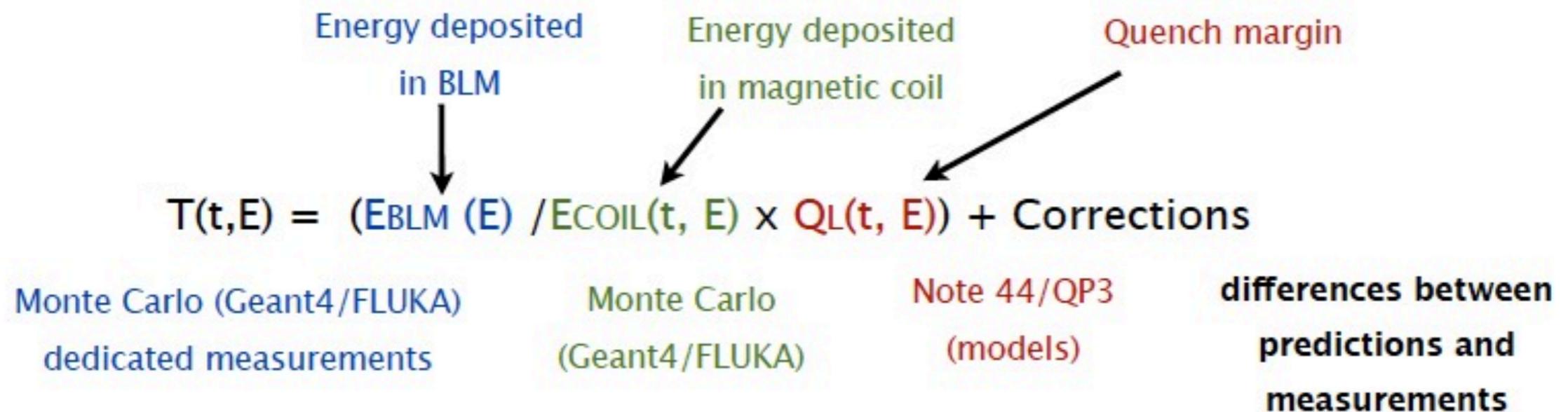
- Avoid quench of superconducting magnets and any damage produced by beam losses.
  - About 3600 Ionization Chambers located around the LHC ring in likely-loss locations
  - The surface electronics cards receive the data for decoding and processing. The system computes the signals integrated in 12 different integration windows. These values are continuously compared to a set of predefined thresholds
  - The surface electronics cards receive the data for decoding and processing. The system computes the signals integrated in 12 different integration windows. These values are continuously compared to a set of predefined thresholds
  - Both the signals recorded by the BLMs and their abort thresholds are sent to Data Bases for offline analysis



# Machine Protection: the LHC case



- In the LHC, the abort thresholds are mainly set to protect against “quench”



- Each BLM has a 12 x 32 threshold table
  - 12 integration windows (40 us, ..., 83 s)
  - 32 energy levels (0 – 7TeV)
- 4000 BLMs x 12 x 32 > 1.5 million thresholds

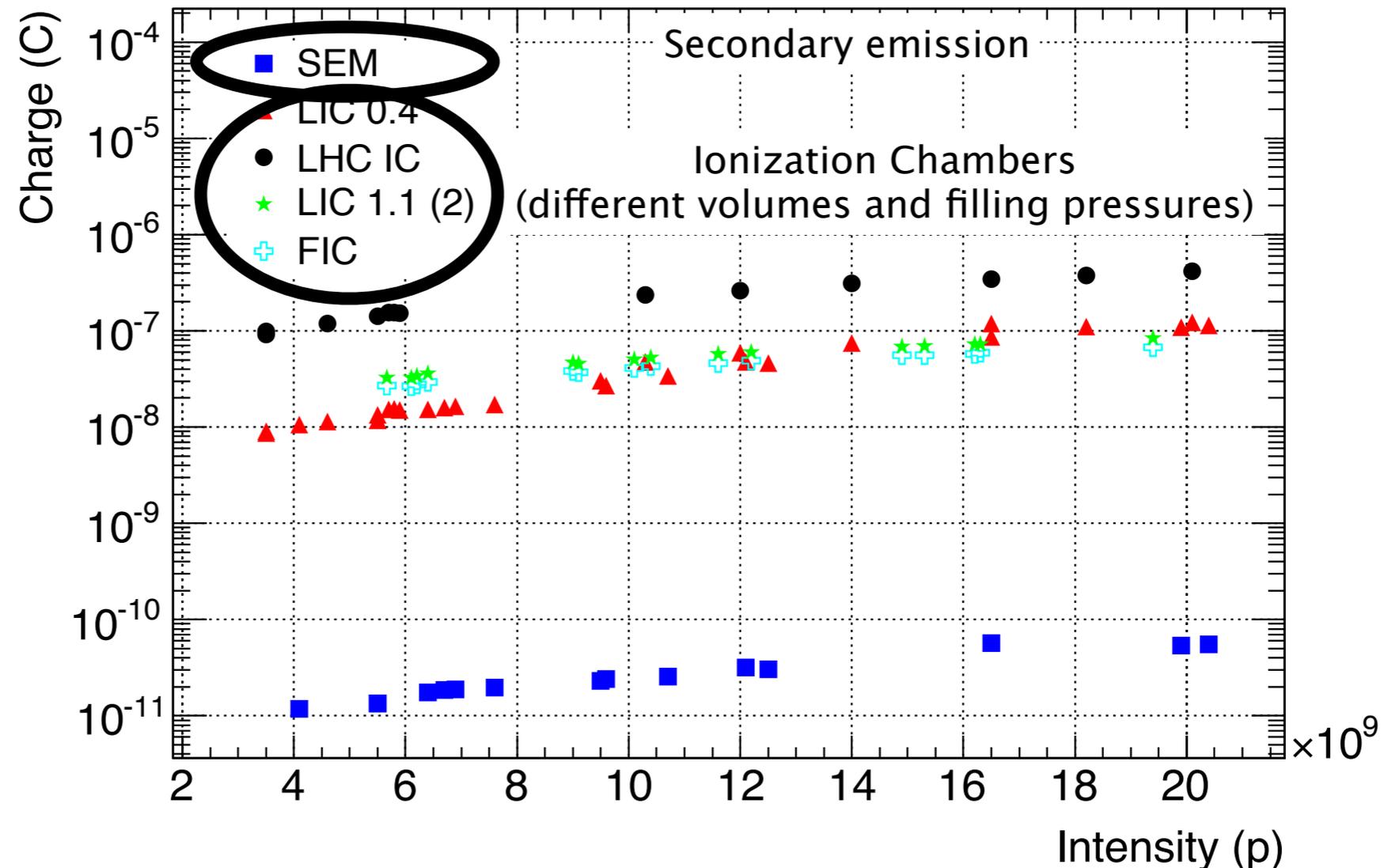


# Challenges

# Dynamic range



- New machines require BLMs covering Dynamic ranges  $\approx 10^{+6}$
- **Various types of BLMs**

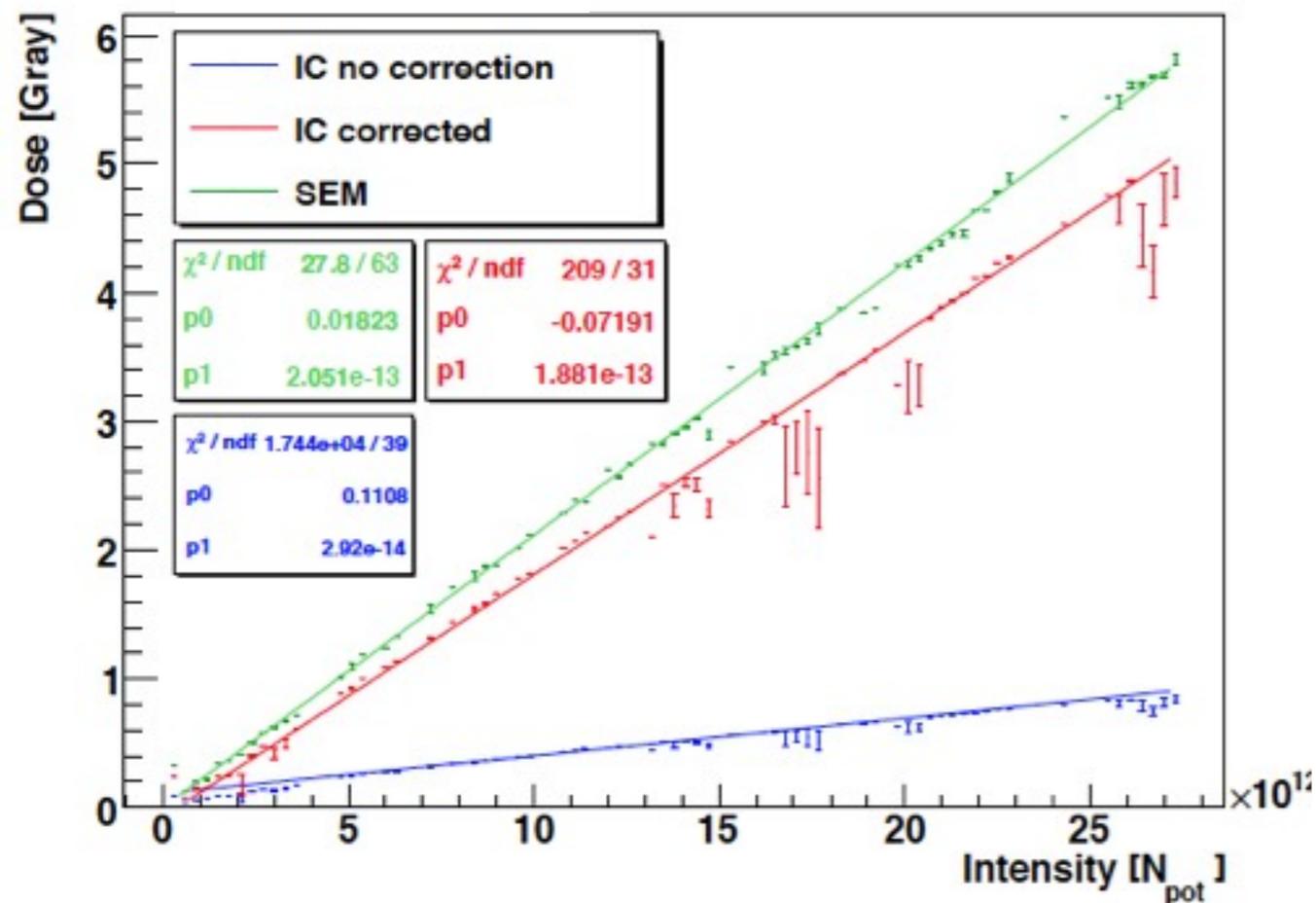


# High radiation environments

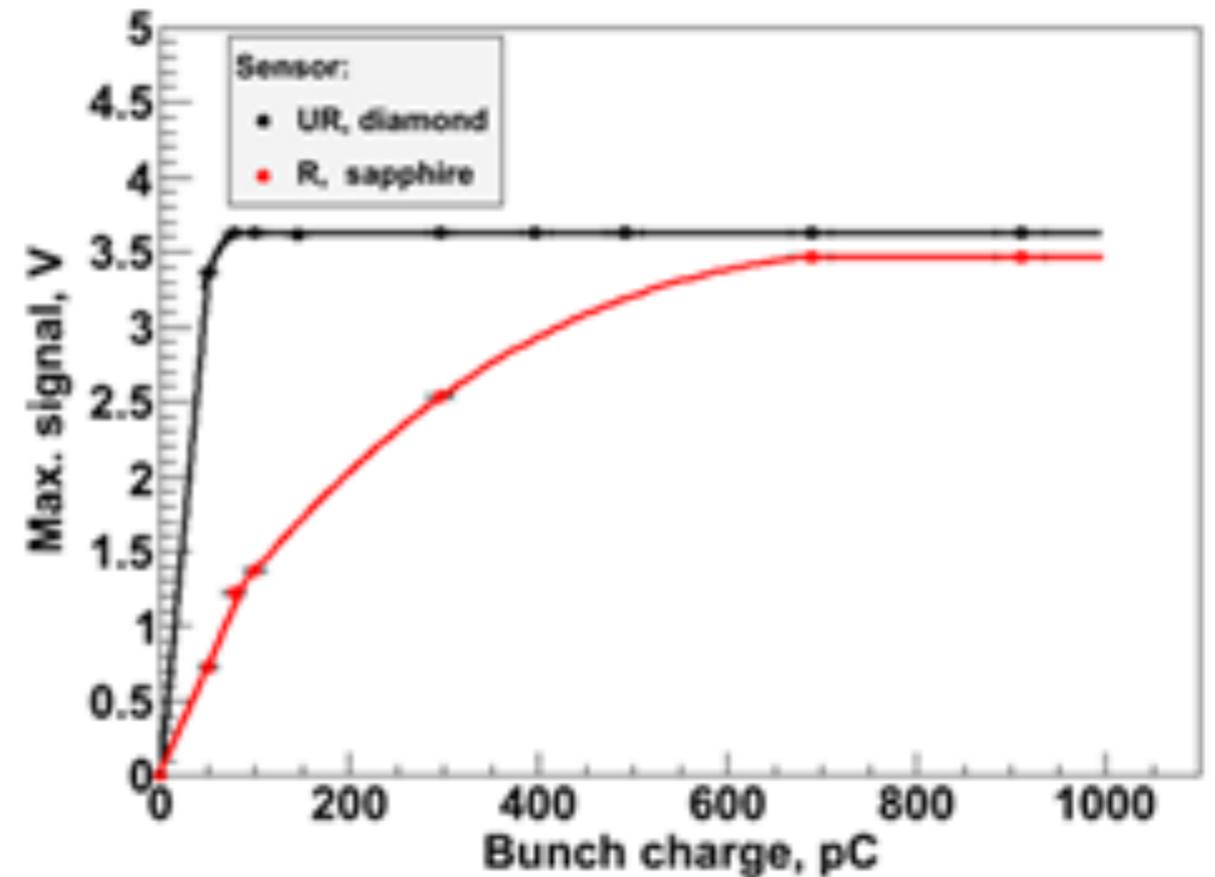


- In hard irradiation conditions most detectors and/or electronics will suffer saturation effects:
  - Understand/correct saturation effects
  - Develop insensitive detectors

D. Kramer Thesis



A. Ignatenko Thesis

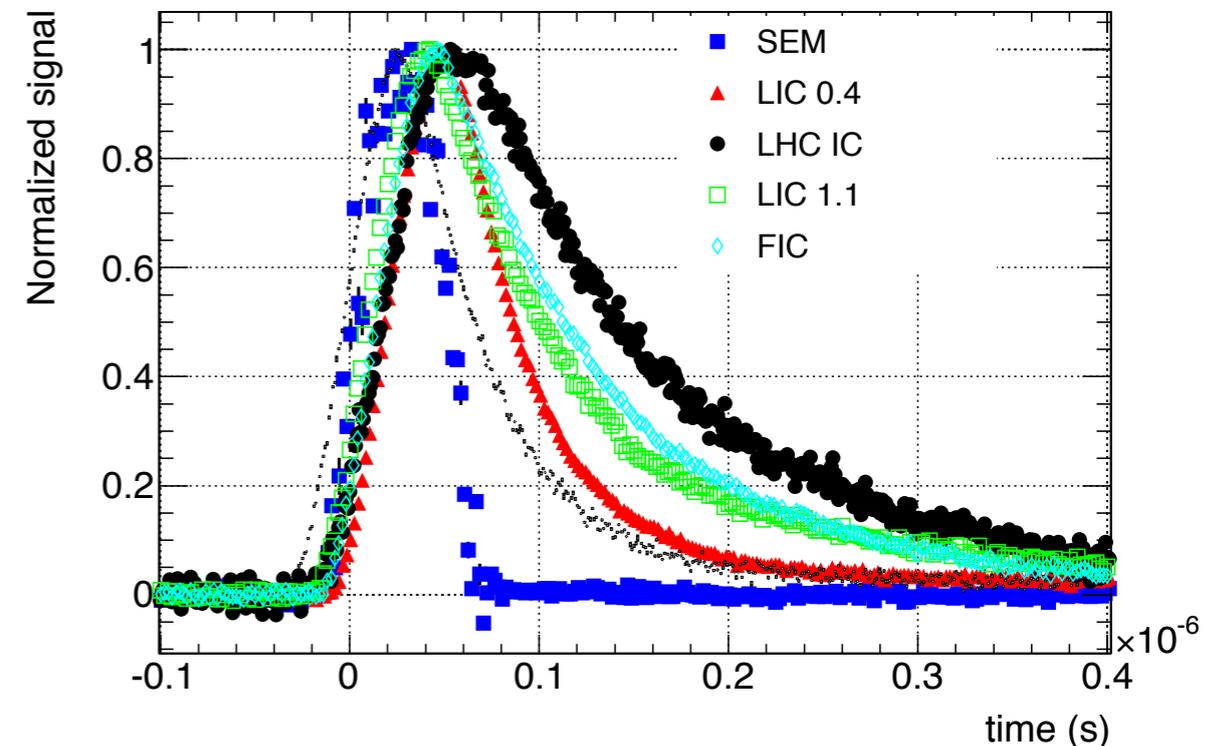


# Time resolution



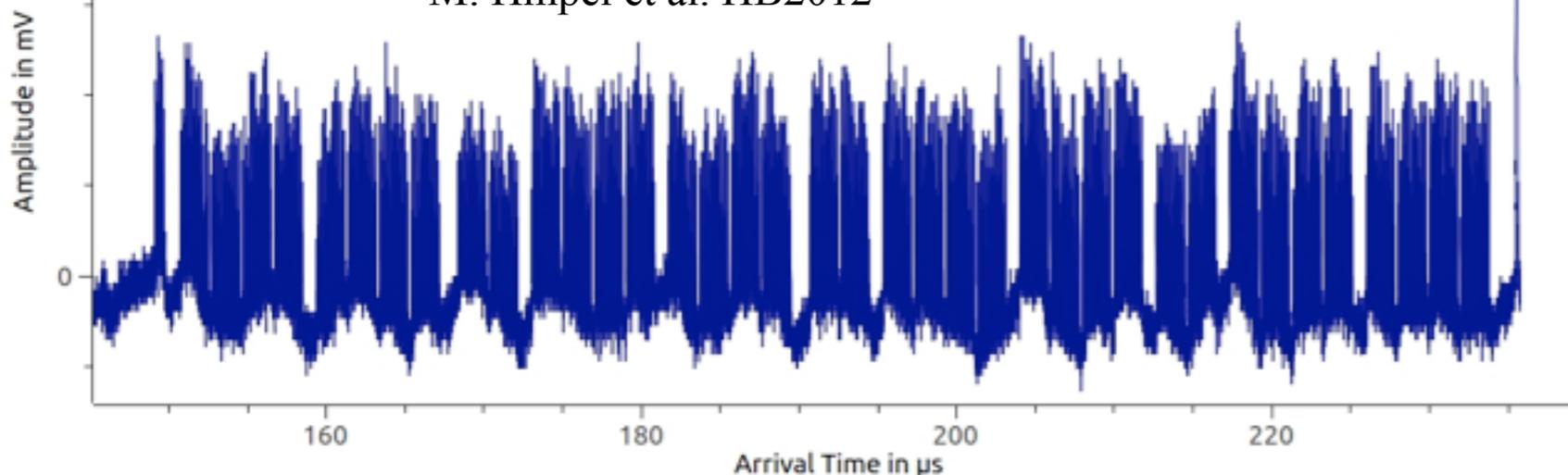
- Determined by the need of measuring bunch by bunch “everything” (that includes beam losses)
- In hadron machines:

- ICs are typical BLMs (radiation hard, reliable, ...)
- But slow ( $\sim \mu\text{s}$ ). Can't see bunch spacing ( $\sim 10\text{--}100\text{ ns}$ )
- SEM are faster but very insensitive. (Broadband amplifiers for  $<10\text{ns}$  resolution BLMs??!!)



## Bunch by bunch beam loss diagnostics with diamond detectors at the LHC

M. Hmpel et al. HB2012



- **Diamond detectors:**
  - ▶ resolution  $\sim 1\text{ns}$
  - ▶ radiation hard

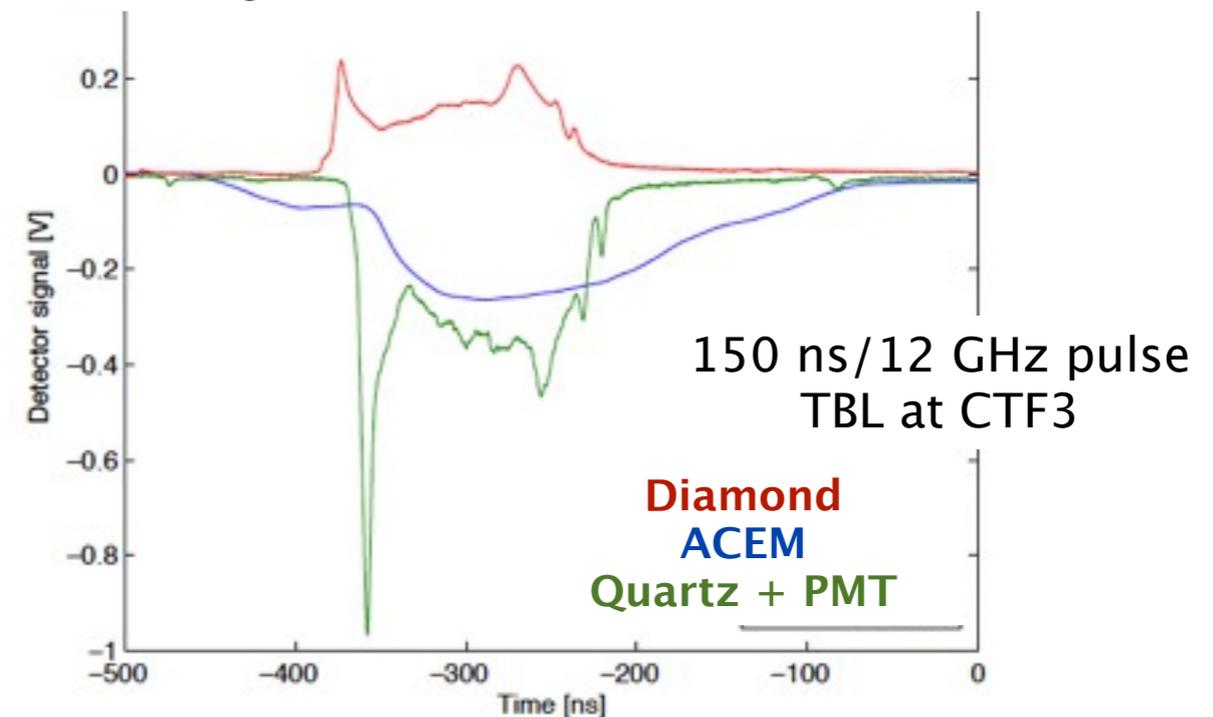
# Time resolution



- In electron machines:
  - The bunch spacing is pushing to unprecedented limits
  - E.g: CLIC: 0.6 – 0.083 ns

## Development of a Beam Loss Monitoring system for CTF3 TBL

E. Branger, Master thesis



- Cherenkov light (as prompt radiation) is currently the best candidate for fast BLMs
  - Ultra fast photo sensor have been demonstrated to disentangle 12GHz bunch spacing

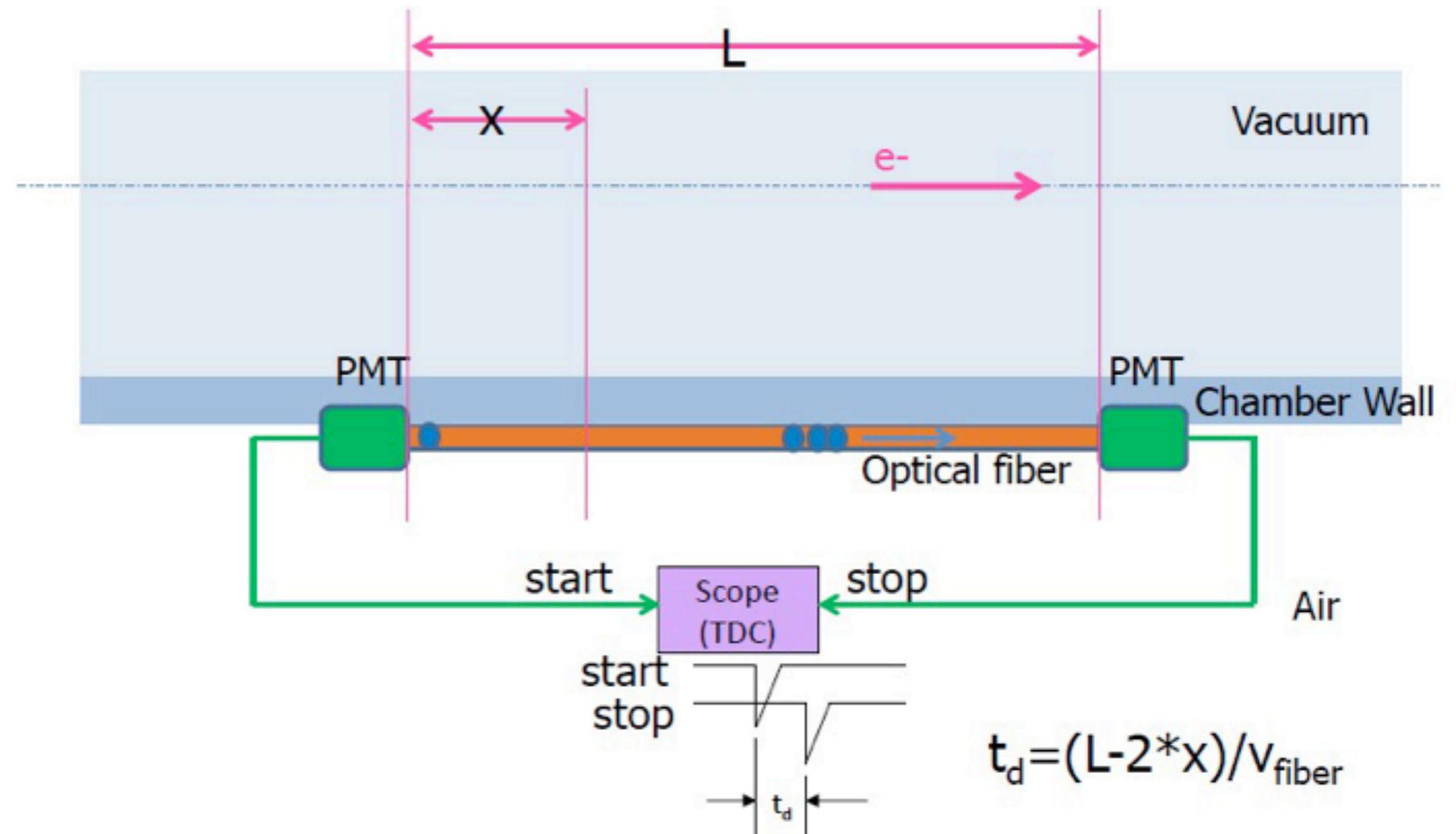
## Application of Metal-Semiconductor-Metal (MSM) Photodetectors for Transverse and Longitudinal Intra-Bunch Beam Diagnostics

R. Steinhagen et al. IBIC13

# Position resolution



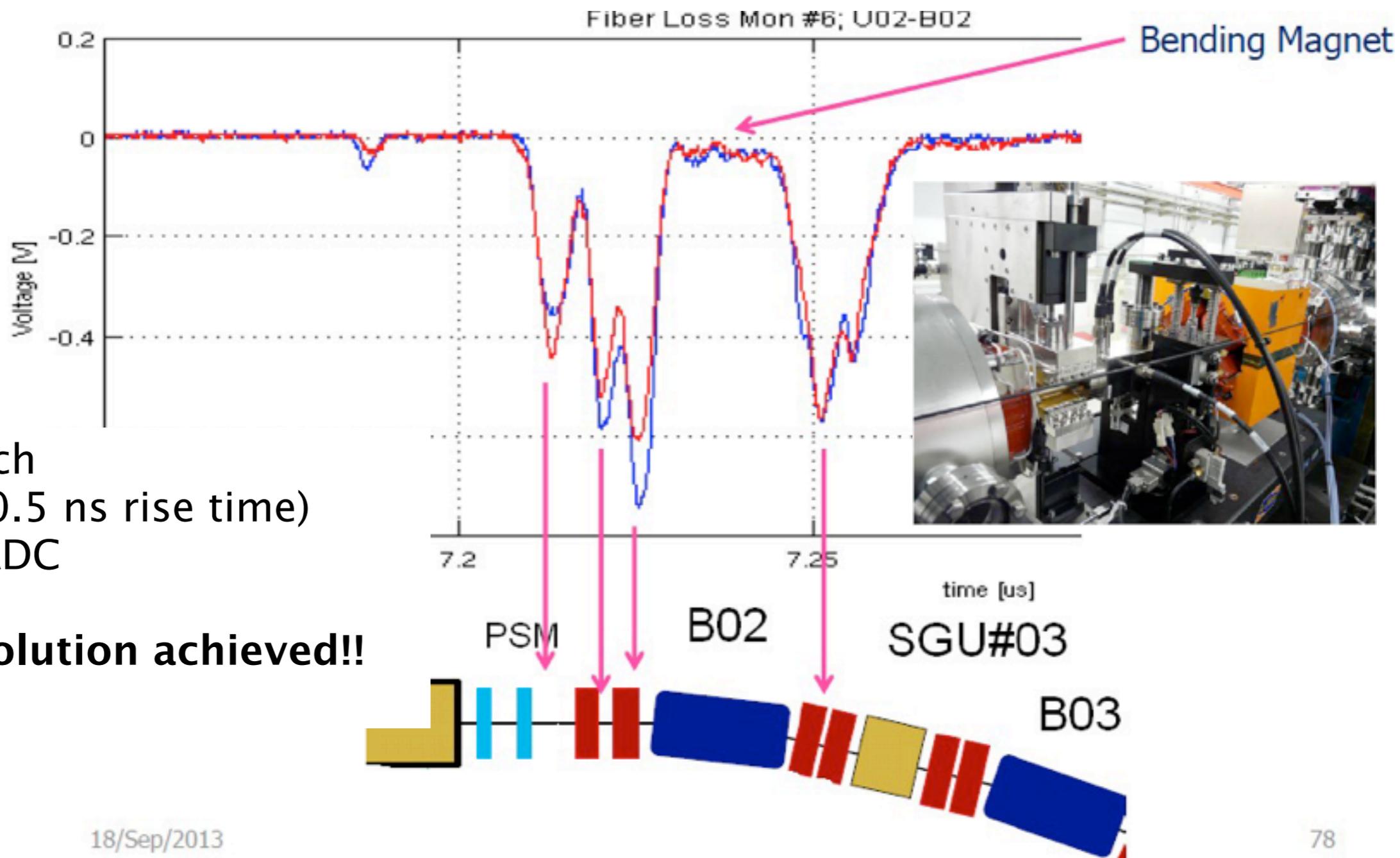
- It is defined by the granularity of the system, i. e. higher number of BLMs provides higher resolution
- Optical fibers provide position resolution by measuring signal delay



# Position resolution



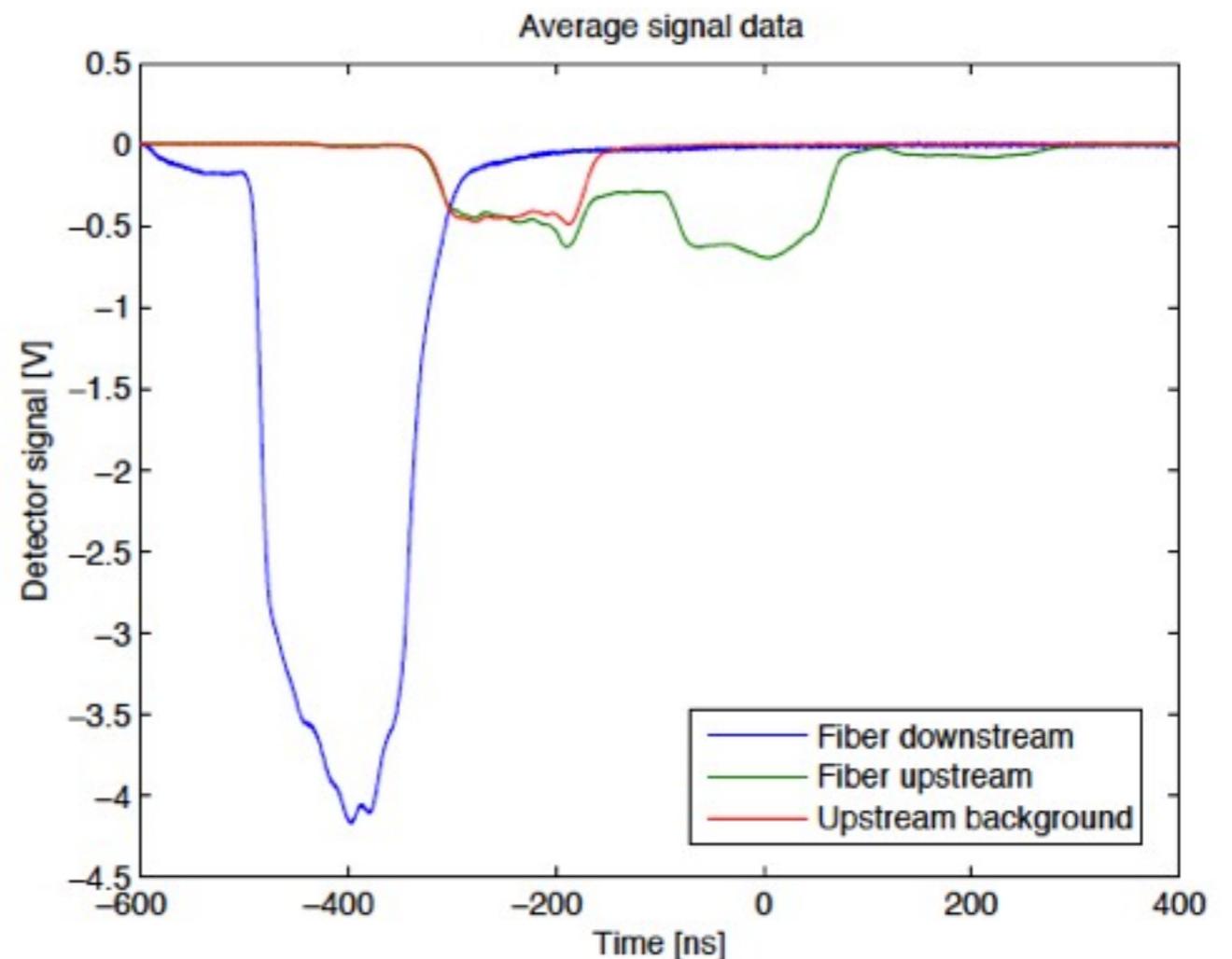
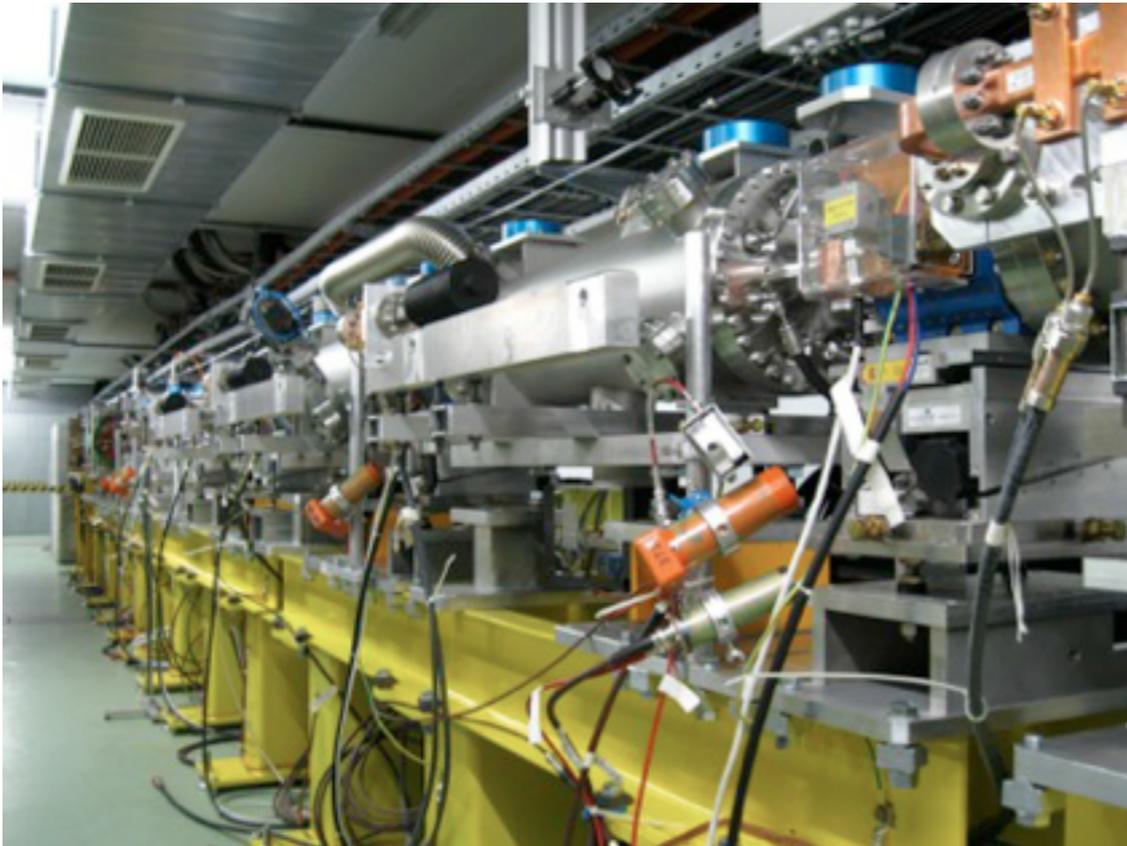
- KEK-PF
- 2.5 GeV electrons
- single bunch



# Position resolution



- What about machines with long pulses?
  - e.g CLIC/CTF3: 150 ns (30 m @2/3c) pulse with 12 GHz bunch spacing
  - Test setup at Test Beam Line (28 m)





- Background = unwanted radiation
  - A BLM system should ideally measure beam losses. However, ionizing radiation can be generated by other mechanisms and detected by your BLMs:
    - ▶ Electromagnetic noise
    - ▶ RF cavities
    - ▶ Synchrotron radiation
  
- Consequence
  - **Dynamic range limitation**

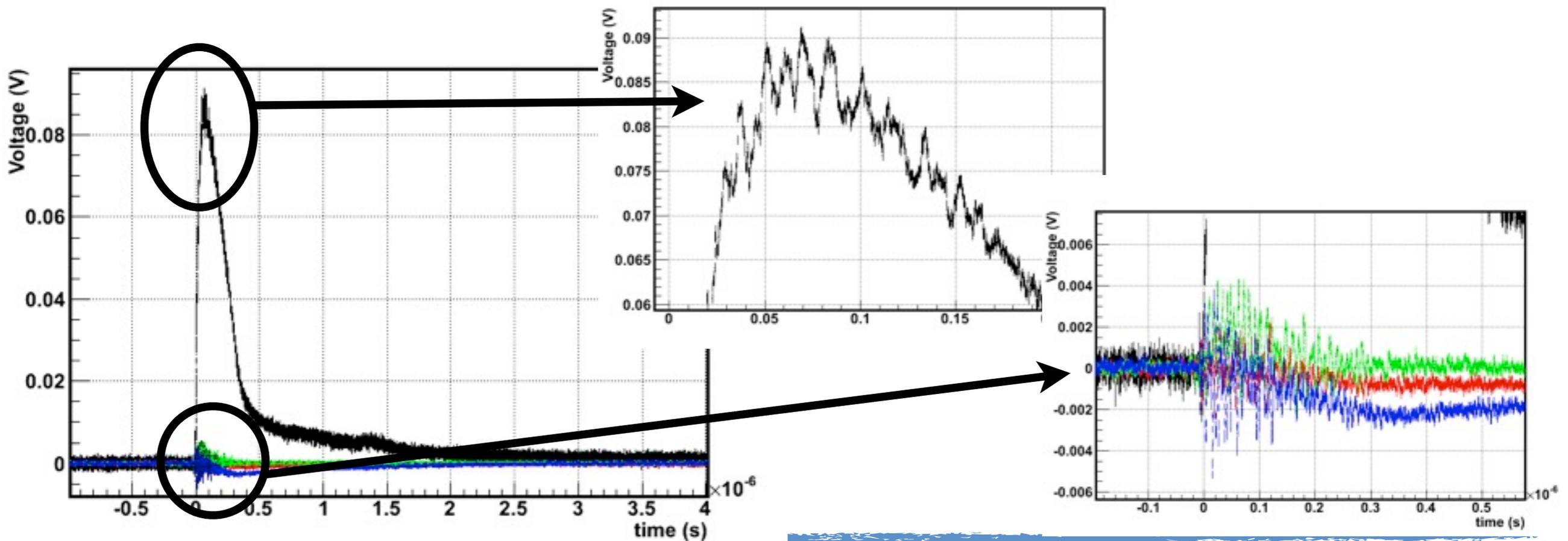
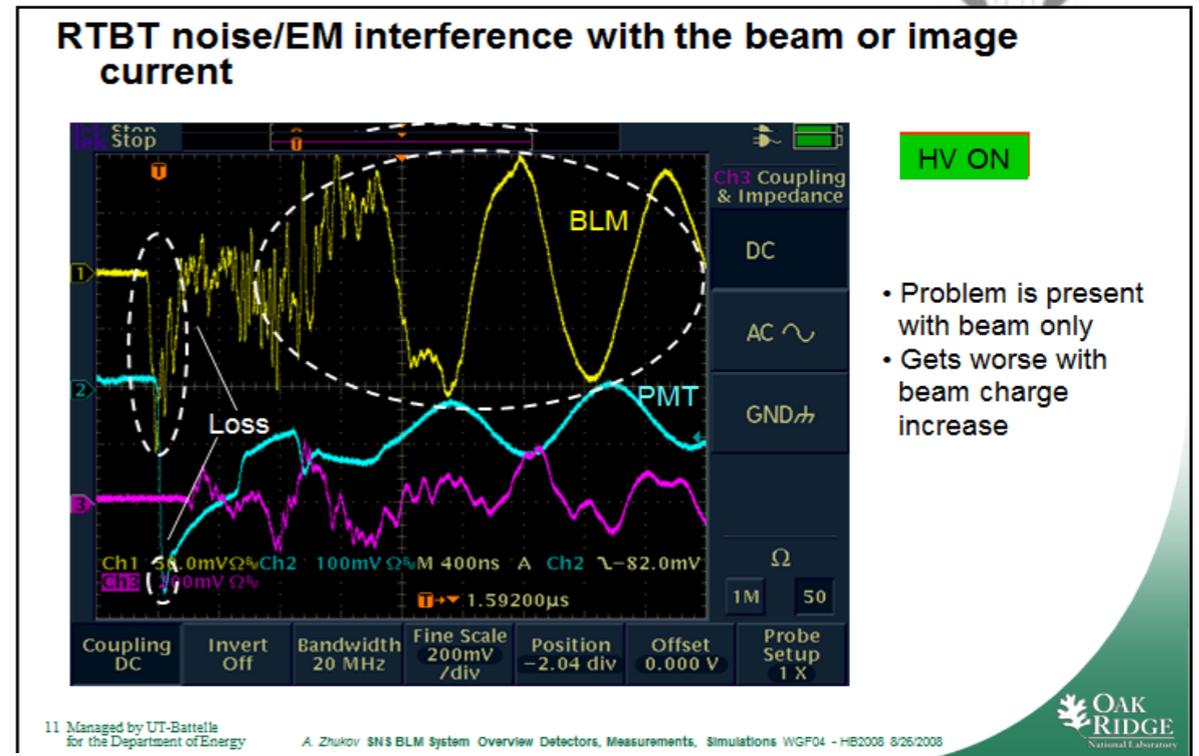
# High background environments



- EM noise

- PS ripple
- Neighboring magnets
- Ground loops

- Provide the best shielding that you can .... and hope for the best!



# High background environments



- RF cavities

- Electrons released (and escaping) in cavity
- X rays

## XFEL Beam Loss Monitor System

A. Kaukher, I. Krouptchenkov et al. IPAC12

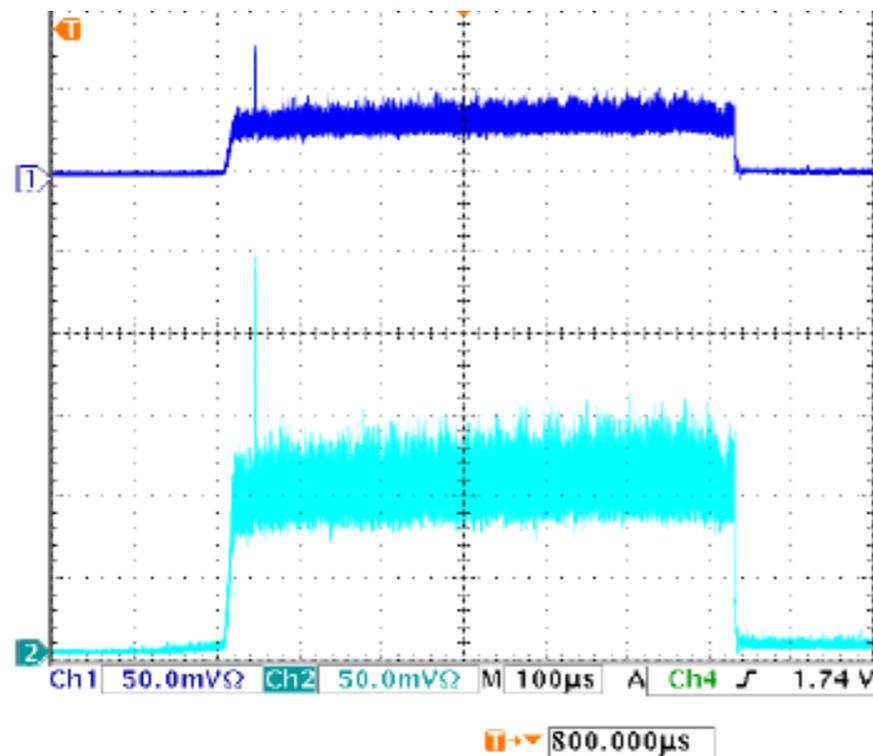


Figure 4: BLM signals from a single bunch and dark current at FLASH (April 2012): BLM with SQ1 synthetic fused silica (top, HV=700 V) and BLM with a scintillator (HV=550 V).

## Beam Loss Detected by Scintillation Monitor

Akihiko Miura, et al. IPAC'11

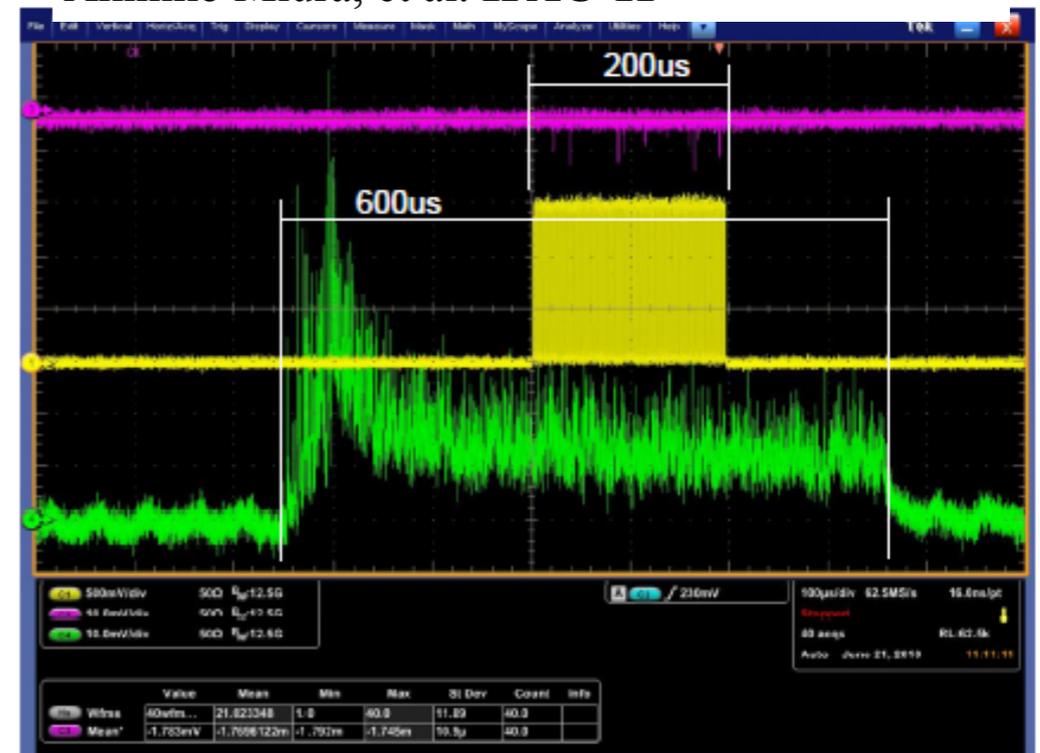


Figure 5: Signals from a gas proportional monitor (green) and plastic scintillation monitor (magenta) at SDTL13 section, during beam operation with chopped beam. The beam current signal with a current transformer is also shown (yellow).

- Subtraction by software (no help on dynamic range) or detectors insensitive to X rays

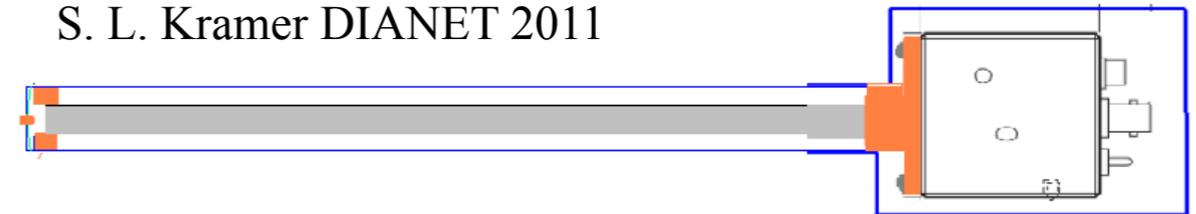
# High background environments



- Synchrotron radiation
  - X rays
- Subtraction by software (no help on dynamic range) or detectors insensitive to X rays

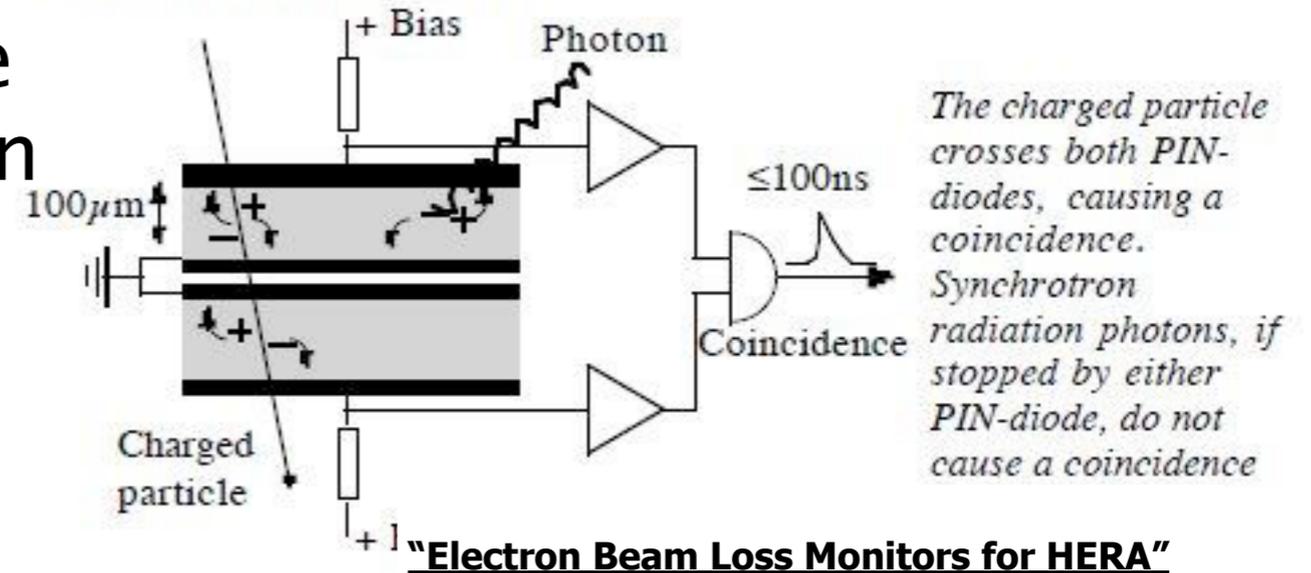
## NSLS-II Beam Loss Monitor System

S. L. Kramer DIANET 2011



- Quartz rods and fibers (as Cherenkov light radiators) are insensitive to neutral radiation
- PIN diode “sandwiches” in coincidence

## Operating principle



“Electron Beam Loss Monitors for HERA”

W. Bialowons et al.

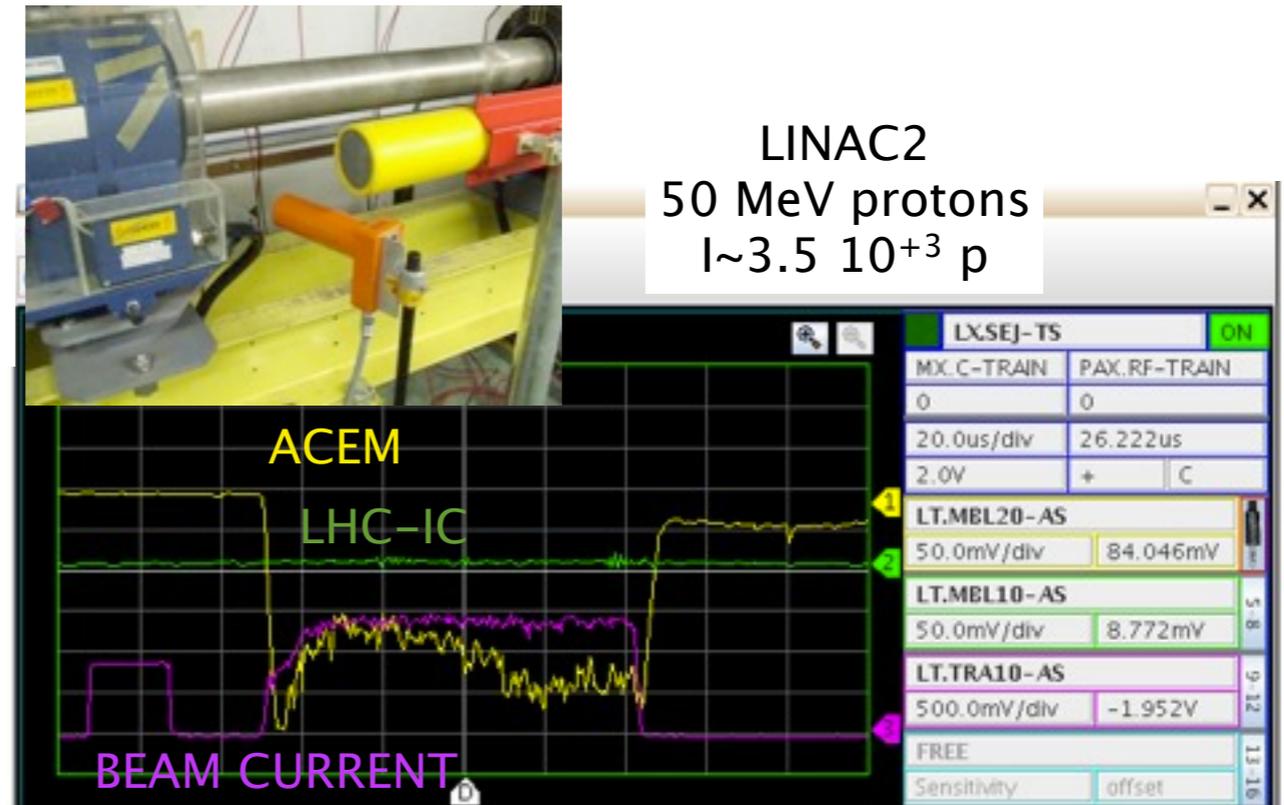
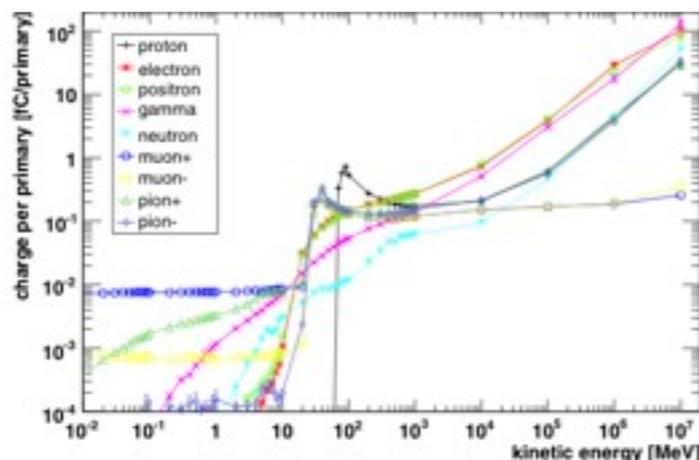
# Beam Losses at low energy



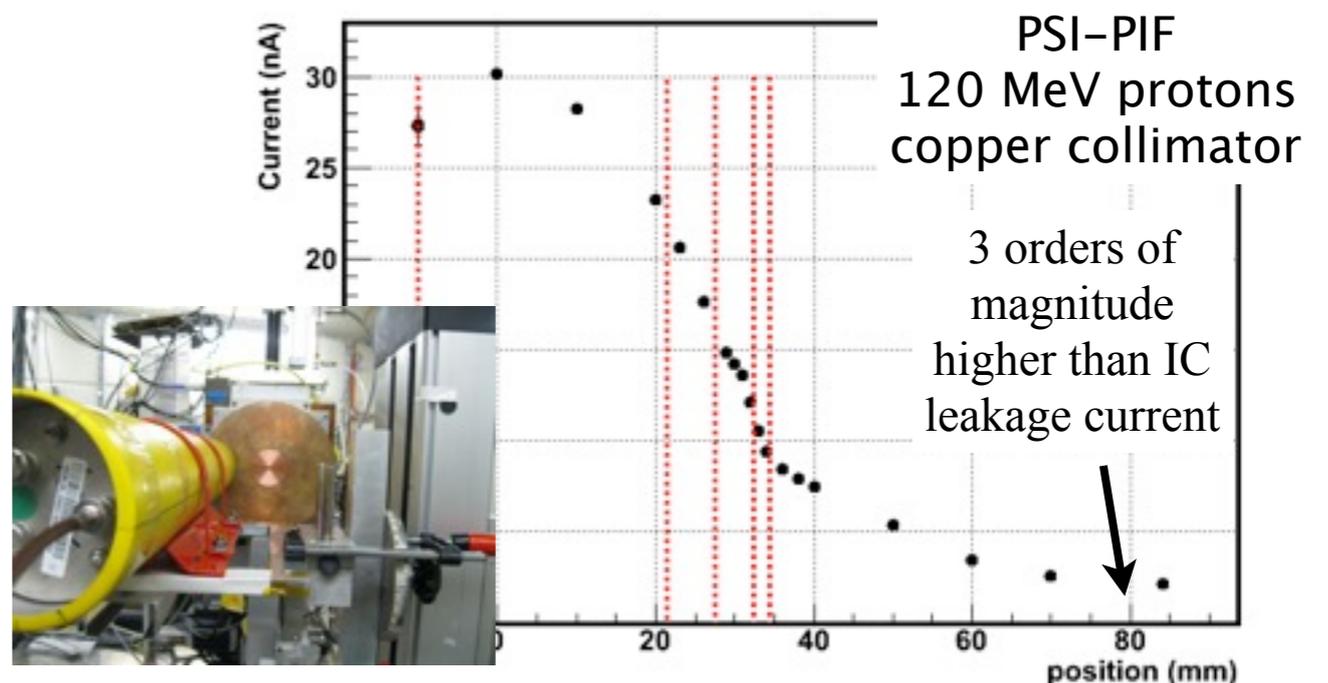
- Very few particles (secondaries) outside of the vacuum chamber
- Even fewer inside the active volume of the detector

## Measurements and Simulations of Ionization Chamber Signals in Mixed Radiation Fields for the LHC BLM System

M. Stockner IEE-NSS 2006



- BLMs very close to beam pipe
- neutron sensitive BLMs
- Very sensitive BLMs (gain)
- BLMs at aperture limitations (collimators)



# Summary and conclusions



- You CANNOT just recycle an already designed BLM system for your own machine
  - Get as much information as possible from past experience
  - Define your specifications according to your accelerator needs
  - Stick with well known technologies if you can (they keep being used for a reason)
- TBL: Implement BLM system within the current operational instrumentation for machine tuning
  - Provide calibrated signals and display
  - Dedicated studies on position resolution with optical fibers



Thank you for your attention

Scintillator material	Density [g/cm <sup>3</sup> ]	Refractive Index	Wavelength [nm] for max. emission	Decay time constant [μs]	Photons/MeV
NaI	3.7	1.78	303	0.06	8 · 10 <sup>4</sup>
NaI(Tl)	3.7	1.85	410	0.25	4 · 10 <sup>4</sup>
CsI(Tl)	4.5	1.80	565	1.0	1.1 · 10 <sup>4</sup>
Bi <sub>4</sub> Ge <sub>3</sub> O <sub>12</sub>	7.1	2.15	480	0.30	2.8 · 10 <sup>3</sup>
CsF	4.1	1.48	390	0.003	2 · 10 <sup>3</sup>
LSO	7.4	1.82	420	0.04	1.4 · 10 <sup>4</sup>
PbWO <sub>4</sub>	8.3	1.82	420	0.006	2 · 10 <sup>2</sup>
LHe	0.1	1.02	390	0.01/1.6	2 · 10 <sup>2</sup>
LAr	1.4	1.29*	150	0.005/0.86	4 · 10 <sup>4</sup>
LXe	3.1	1.60*	150	0.003/0.02	4 · 10 <sup>4</sup>

\* at 170 nm

Scintillator material	Density [g/cm <sup>3</sup> ]	Refractive Index	Wavelength [nm] for max. emission	Decay time constant [ns]	Photons/MeV
Naphtalene	1.15	1.58	348	11	$4 \cdot 10^3$
Antracene	1.25	1.59	448	30	$4 \cdot 10^4$
p-Terphenyl	1.23	1.65	391	6-12	$1.2 \cdot 10^4$
NE102*	1.03	1.58	425	2.5	$2.5 \cdot 10^4$
NE104*	1.03	1.58	405	1.8	$2.4 \cdot 10^4$
NE110*	1.03	1.58	437	3.3	$2.4 \cdot 10^4$
NE111*	1.03	1.58	370	1.7	$2.3 \cdot 10^4$
BC400**	1.03	1.58	423	2.4	$2.5 \cdot 10^2$
BC428**	1.03	1.58	480	12.5	$2.2 \cdot 10^4$
BC443**	1.05	1.58	425	2.2	$2.4 \cdot 10^4$

Sensitivity  $S_{Ion}$  [C/rad] of Ion-chamber depends on geometry!

Example: 1 ltr. Argon filled chamber, 100% charge sampling efficiency:

$$\rho = 1.661 \cdot 10^{-3} \text{ g/cm}^3 \text{ (20}^{\circ}\text{C, 1 atm)}$$

$$S_{Ion} = \frac{100 \frac{\text{erg}}{\text{g}}}{1 \text{ rad}} \cdot \frac{1 \text{ eV}}{1.6 \cdot 10^{-12} \text{ erg}} \cdot \frac{1 \text{ e}^-}{26 \text{ eV}} \cdot \frac{1.66 \cdot 10^{-3} \text{ g}}{\text{cm}^3} \cdot 1000 \text{ cm}^3 \cdot \frac{1.6 \cdot 10^{-19} \text{ C}}{\text{e}^-} = 638 \frac{\text{nC}}{\text{rad}}$$

1 rad
eV → erg
W
ρ
1 ltr.
e<sup>-</sup> charge



### Short Ionization Chamber



### Calibration/Sensitivity

The number of electrons ( $n_e$ ) produced in the gap by one minimum ionizing Particles (MIPs) is:

$$n_e = \frac{D \cdot \rho}{W} \cdot \frac{dE}{dx} (\text{Medium}) \leftarrow \text{Bethe-Bloch}$$

Note that the average energy needed to produce an electron-ion pair (W-factor) is larger than the ionization energy. Is about constant for many gases and radiations.

Example: Argon:

$$\rho = 1.661 \cdot 10^{-3} \text{ g/cm}^3 \text{ (20}^{\circ}\text{C, 1 atm)}$$

$$dE/dx = 1.52 \text{ MeV/(g/cm}^2\text{)}$$

Gas	first ionisation potential	fast electrons
Ar	15.7	26.4
He	24.5	41.3
H <sub>2</sub>	15.6	36.5
N <sub>2</sub>	15.5	34.8
Air		33.8
O <sub>2</sub>	12.5	30.8
CH <sub>4</sub>	14.5	27.3

$$n_e \approx 100/\text{cm} \cdot D \text{ [e/ MIP]}$$

Note: Cross section for nuclear interaction is about  $5 \cdot 10^{-6}$  times the ionization cross section ( $10^{-16} \text{ cm}^2$ ). Cross section for excitation is about  $10^{-17} \text{ cm}^2$ . Rutherford (nuclear) scattering does not produce significant energy transfer, but angular spread.